

Giant elastoresistance in magic-angle twisted bilayer graphene

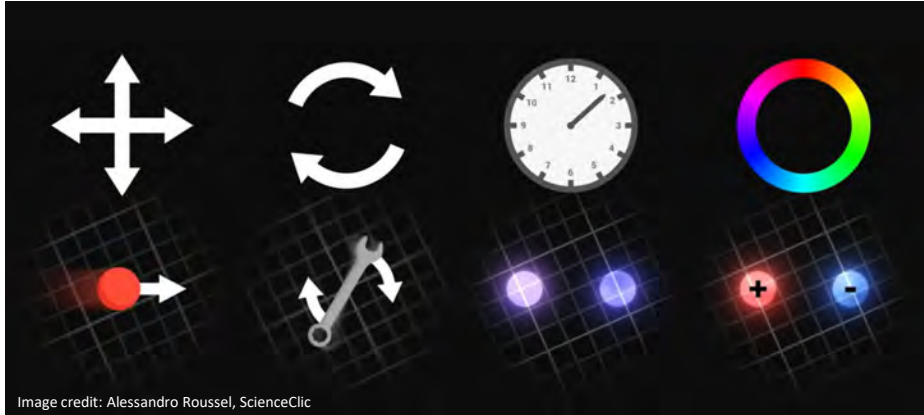
Matthew Yankowitz

Physics, MSE
University of Washington

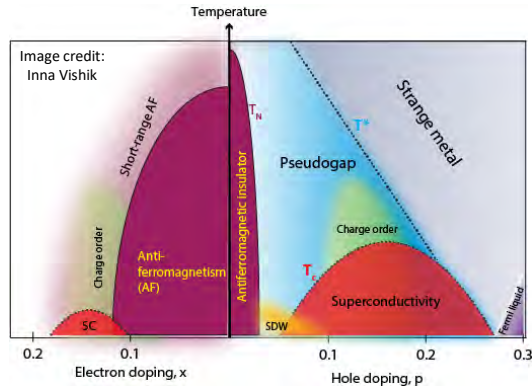
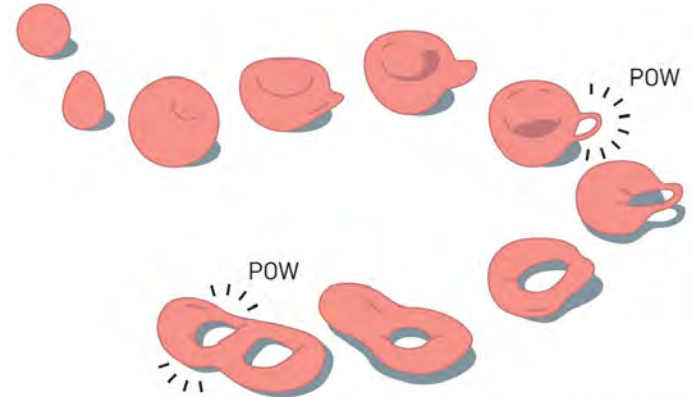


Organizing principles for condensed matter physics

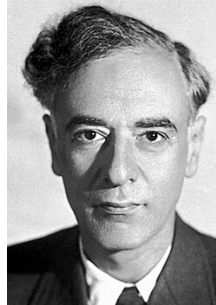
Symmetry



Topology



Lev Landau

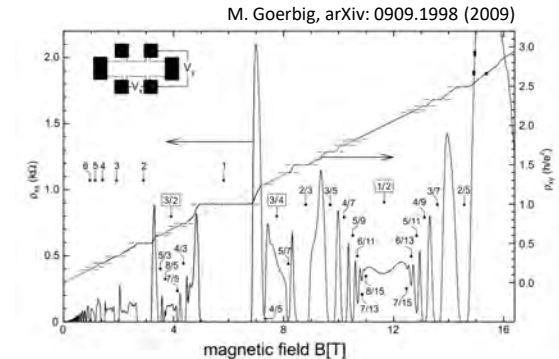


phases described by their broken symmetries

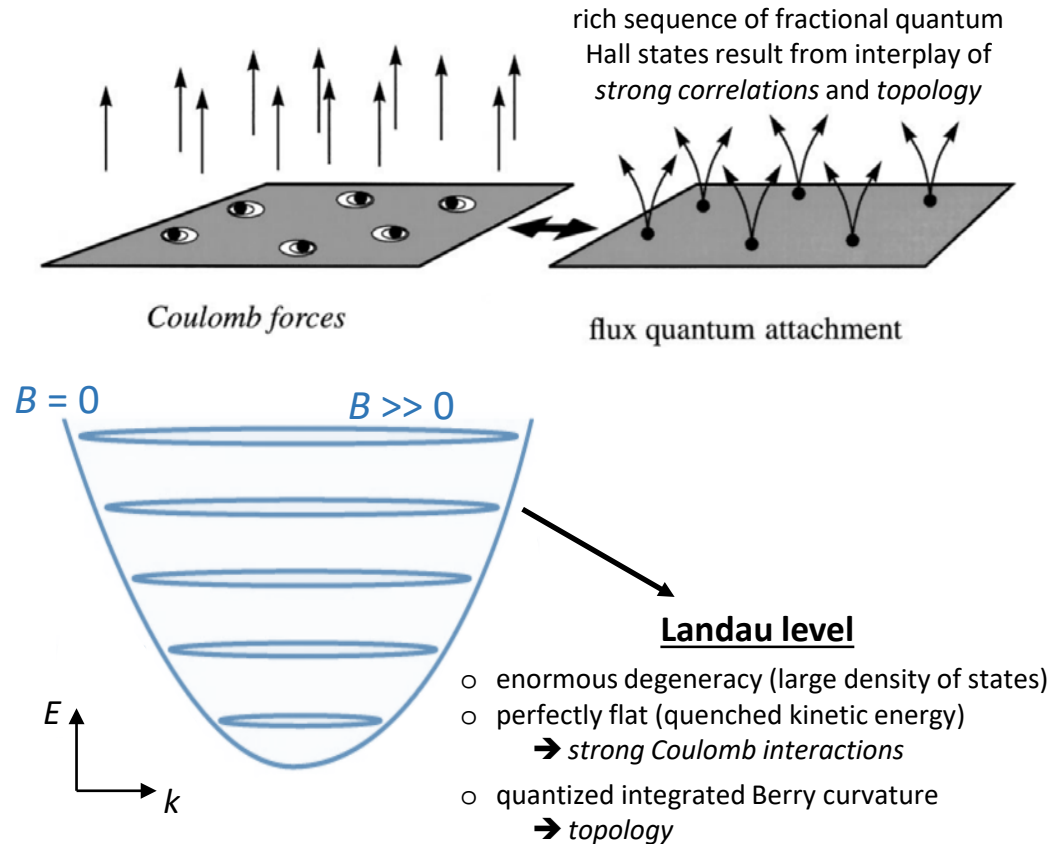
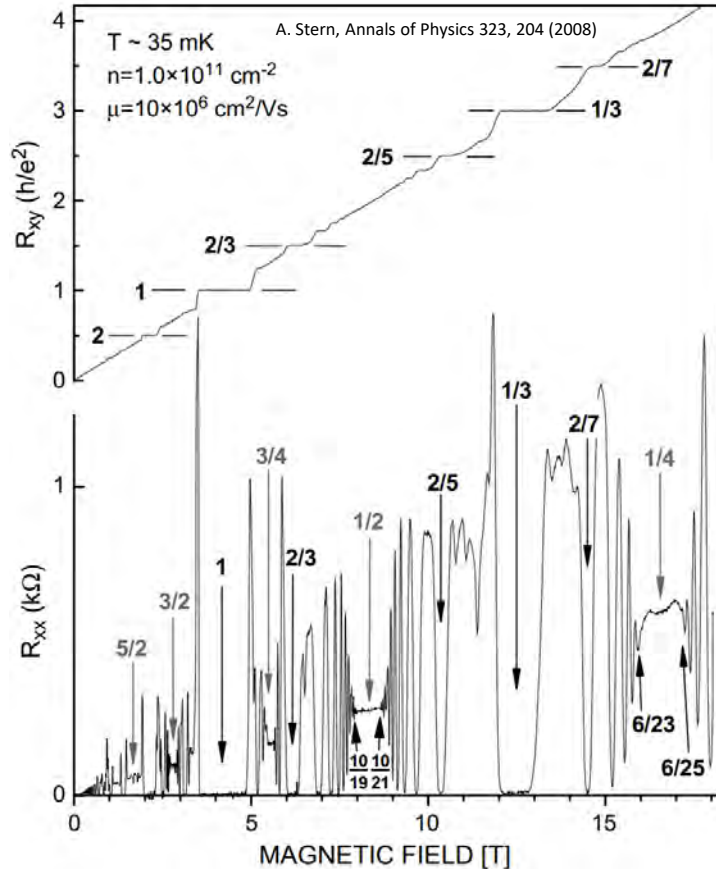
David Thouless



phases described by their topological invariants



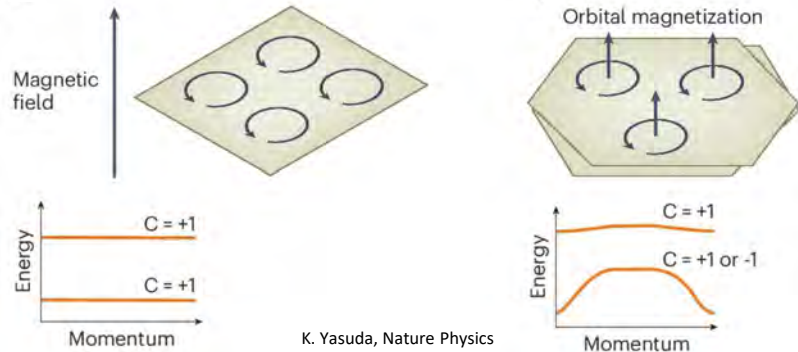
Intertwining of strong correlations and topology



Where the Topological Flat Bands Are

Need to have:

- flat electronic dispersion
(large ratio of Coulomb-to-kinetic energy)
- Berry curvature
(effective magnetic field from atomic arrangement)



K. Yasuda, Nature Physics
21, 499 (2025)



Options:

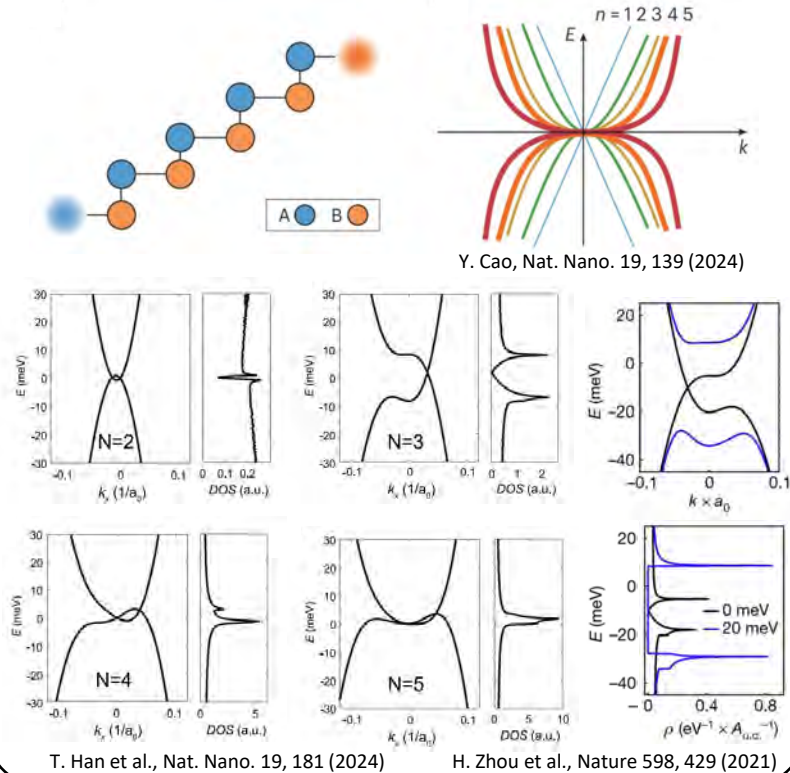
- *discover* a new material (...a “gift from nature”)
- *engineer* a new material (...do it ourselves!)

Graphite: all we need

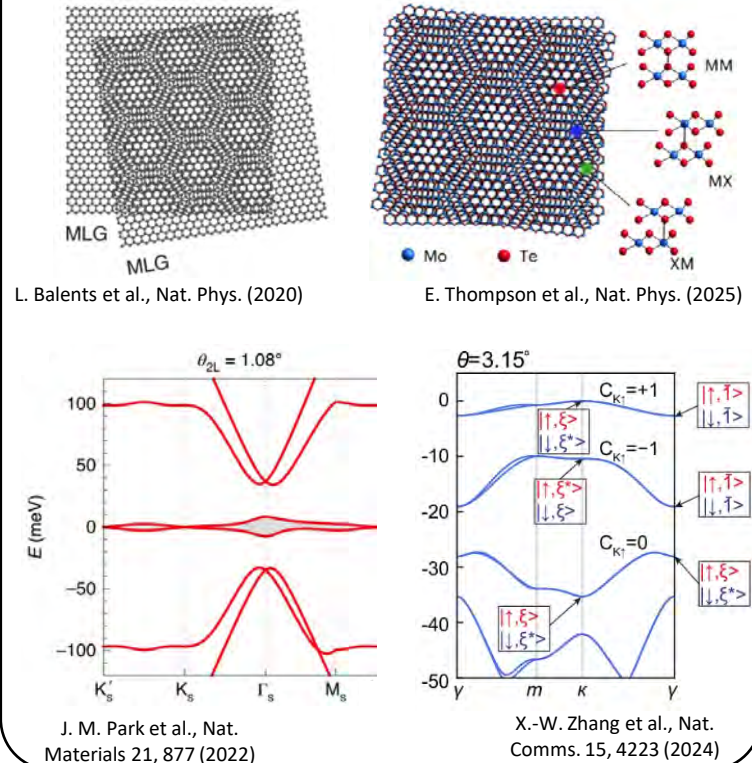


Designing topological flat bands in 2D materials

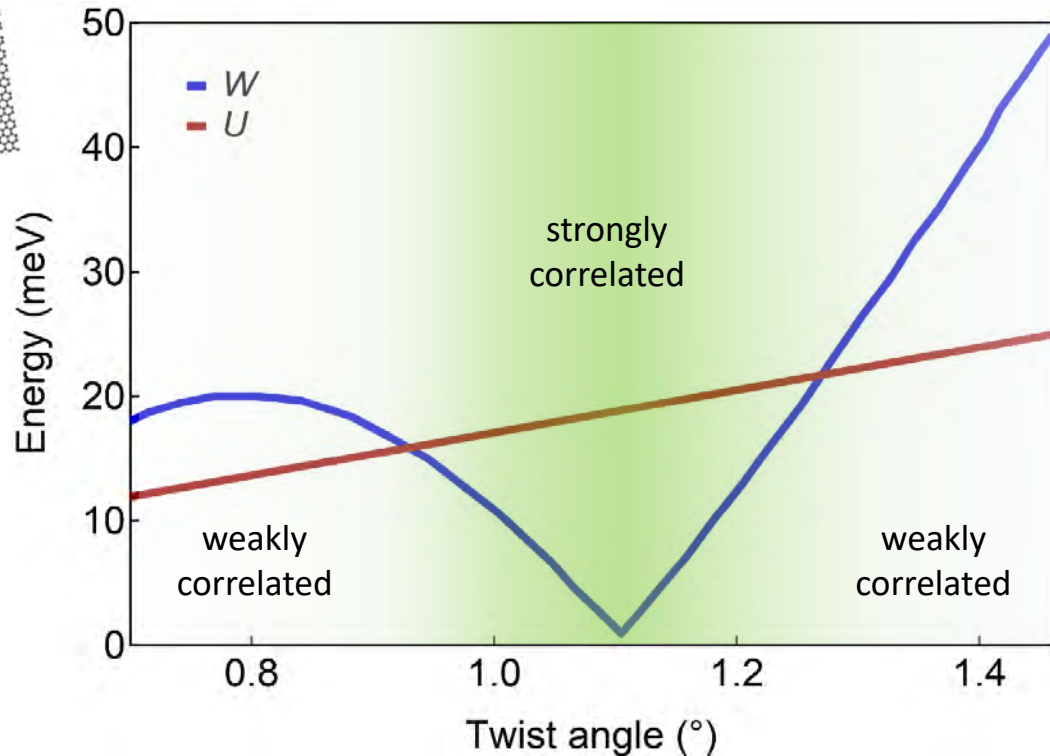
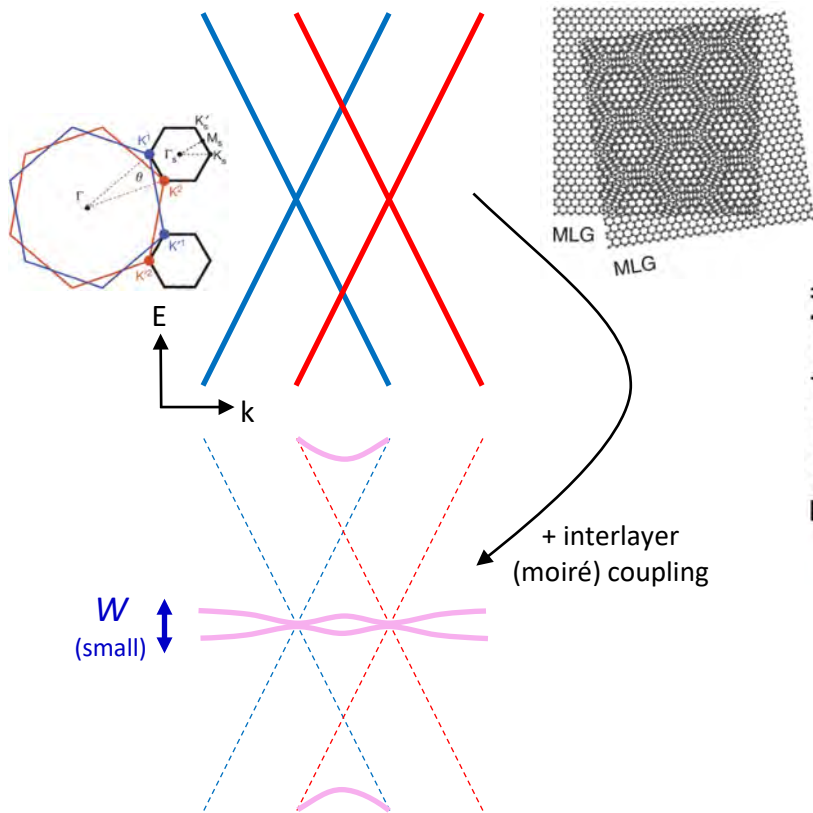
crystalline multilayers



moiré multilayers

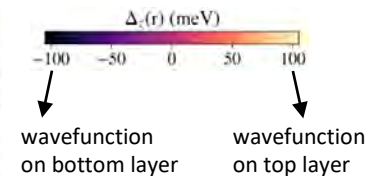
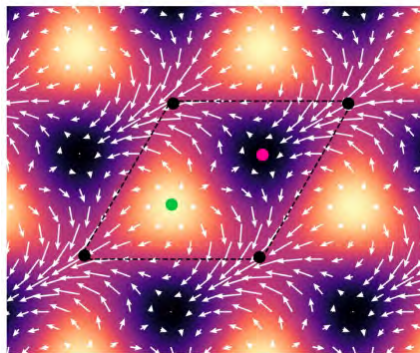
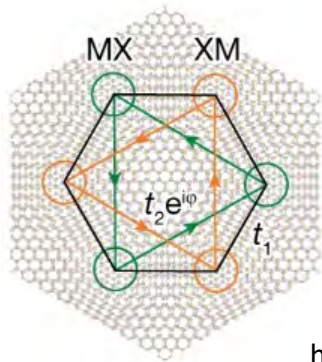


Band flattening in twisted bilayer graphene



Topology from Berry curvature in moiré materials

Kane-Mele model



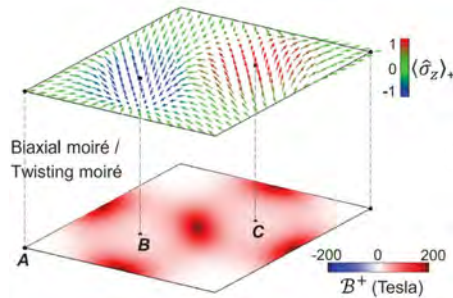
X. Zhang, arXiv:2311.12776 (2024)
See also: F. Wu, PRL 122, 086402 (2019)

real-space Berry curvature analog:

$$\dot{\mathbf{k}} = -\frac{\partial E}{\partial \mathbf{R}} + \dot{\mathbf{R}} \times \boldsymbol{\Omega}^R$$

$$\boldsymbol{\Omega}^R \equiv i \left\langle \frac{\partial u}{\partial \mathbf{R}} \left| \times \right| \frac{\partial u}{\partial \mathbf{R}} \right\rangle$$

superlattice-generated
effective magnetic field



H. Yu, National Science Review 7, 12 (2020)

how does moiré superlattice set band topology?

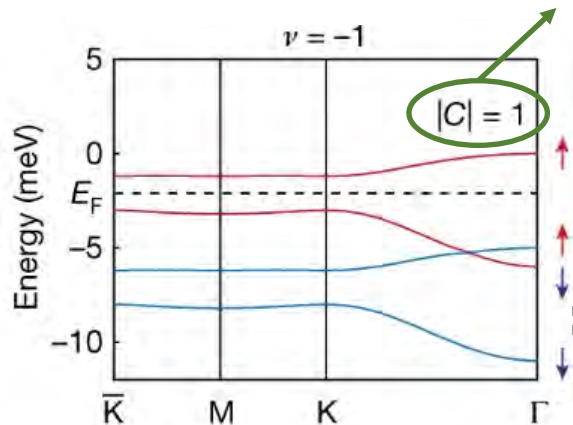
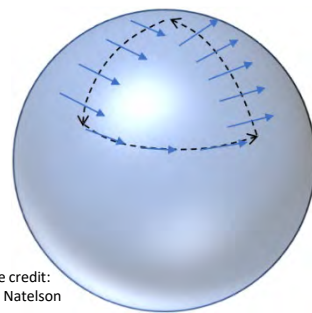


Image credit:
Doug Natelson



...through its Berry phases

H. Park, Nature 622, 74 (2023)

"Brute force" calculation:

$$\dot{\mathbf{R}} = \frac{\partial E}{\partial \mathbf{k}} - \dot{\mathbf{k}} \times \boldsymbol{\Omega}^k$$

(anomalous velocity)

$$\boldsymbol{\Omega}^k \equiv i \left\langle \frac{\partial u}{\partial \mathbf{k}} \left| \times \right| \frac{\partial u}{\partial \mathbf{k}} \right\rangle$$

(Berry curvature)

$$\mathcal{C} = \frac{1}{2\pi} \mathbf{e}_z \cdot \int_{\text{mBZ}} \boldsymbol{\Omega}^k d\mathbf{k}$$

Microscopic fingerprints of topology



Ellis Thompson

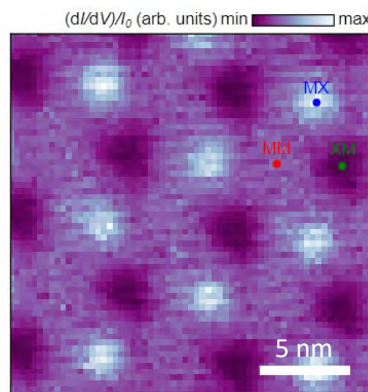
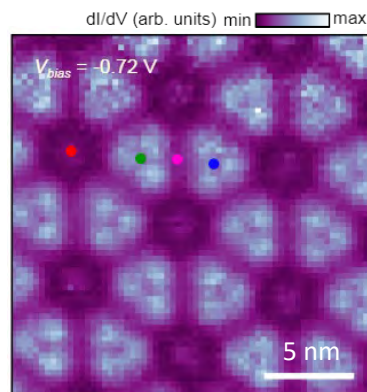
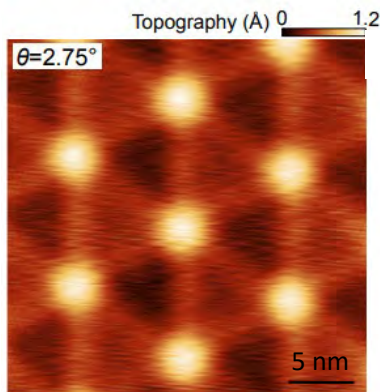
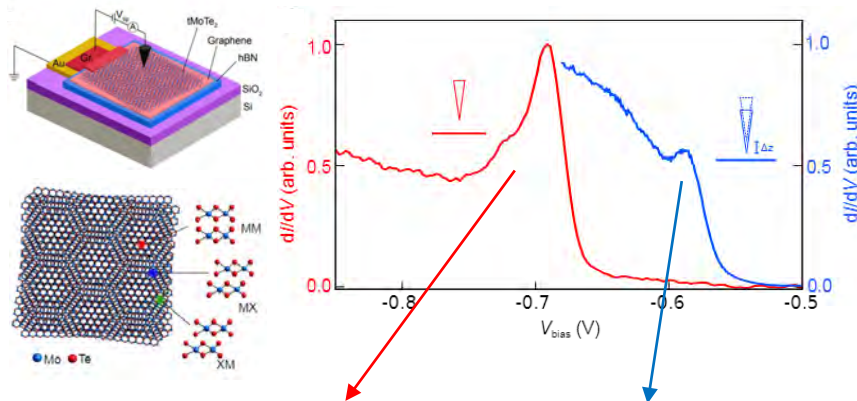


Toby Chu



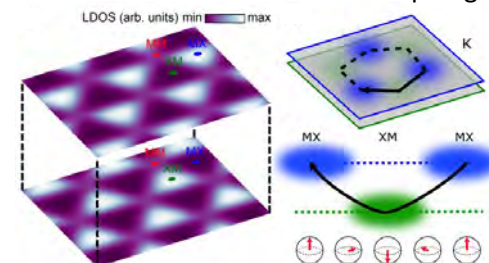
Dr. Florie Mesple

Prof. Xiaodong Xu
Prof. Ting Cao
Prof. Di Xiao
Dr. Xiao-Wei Zhang

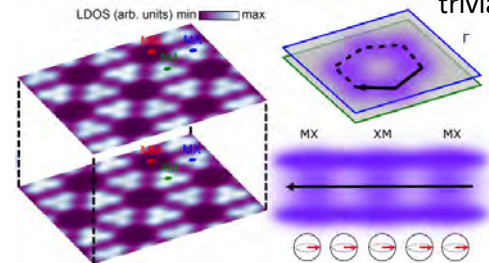


DFT calculation

topological ($C = 1$)



trivial ($C = 0$)



Robustness of real-space topology in moiré systems

Krystof Kolár,^{1,2} Kang Yang,¹ Felix von Oppen,¹ and Christophe Mora³

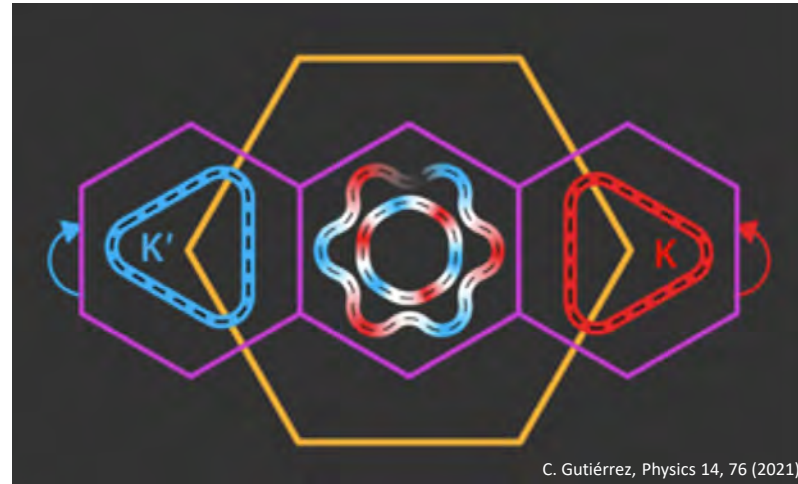
¹Dahlem Center for Complex Quantum Systems and Fachbereich Physik, Freie Universität Berlin, 14195 Berlin, Germany

²Department of Applied Physics, Aalto University School of Science, FI-00076 Aalto, Finland

³Université Paris Cité, CNRS, Laboratoire Matériaux et Phénomènes Quantiques, 75013 Paris, France
(Dated: July 2, 2025)

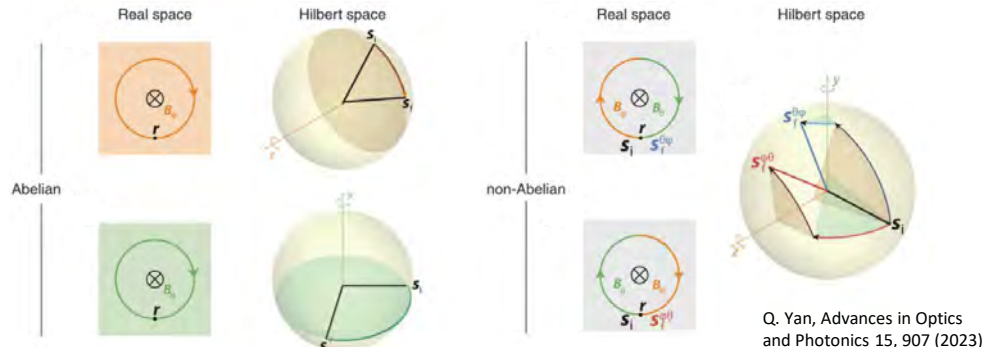
Graphene platforms for new quantum effects

flat bands
→ *strong correlations*
and *symmetry breaking*

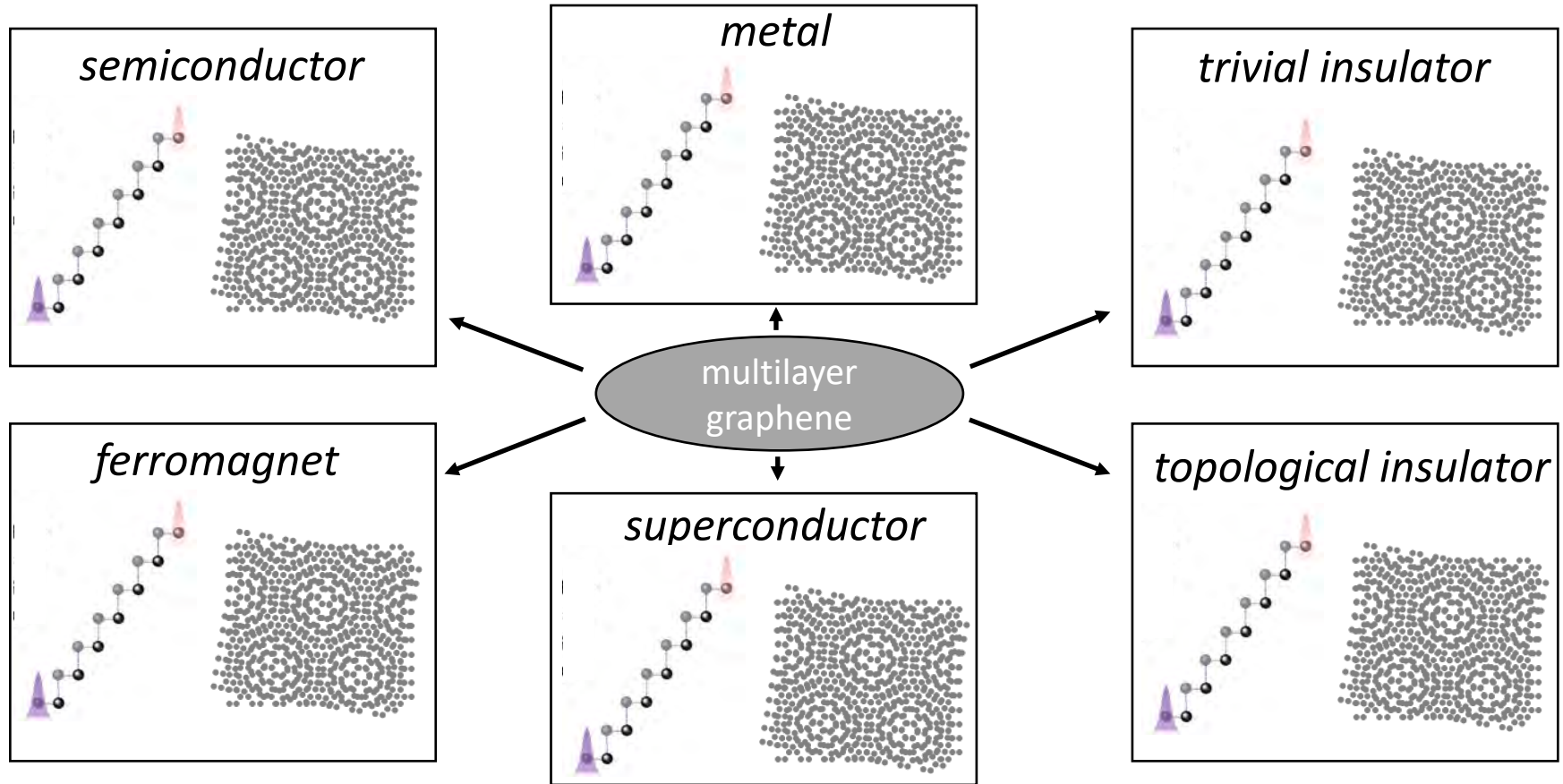


C. Gutiérrez, Physics 14, 76 (2021)

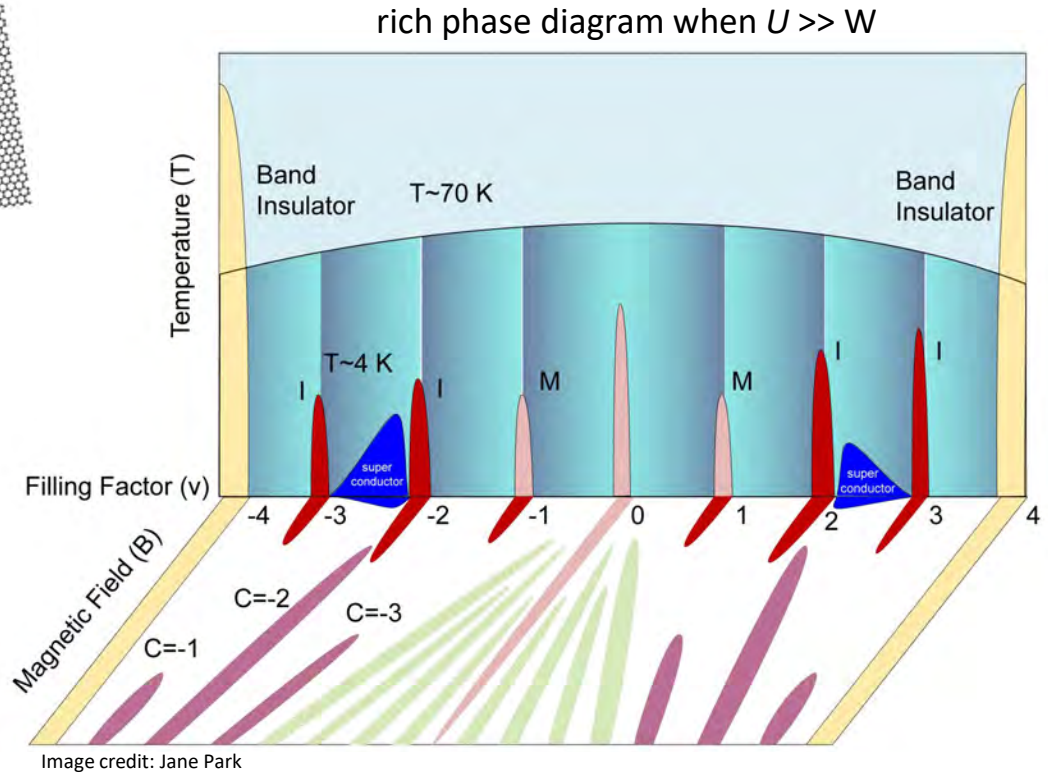
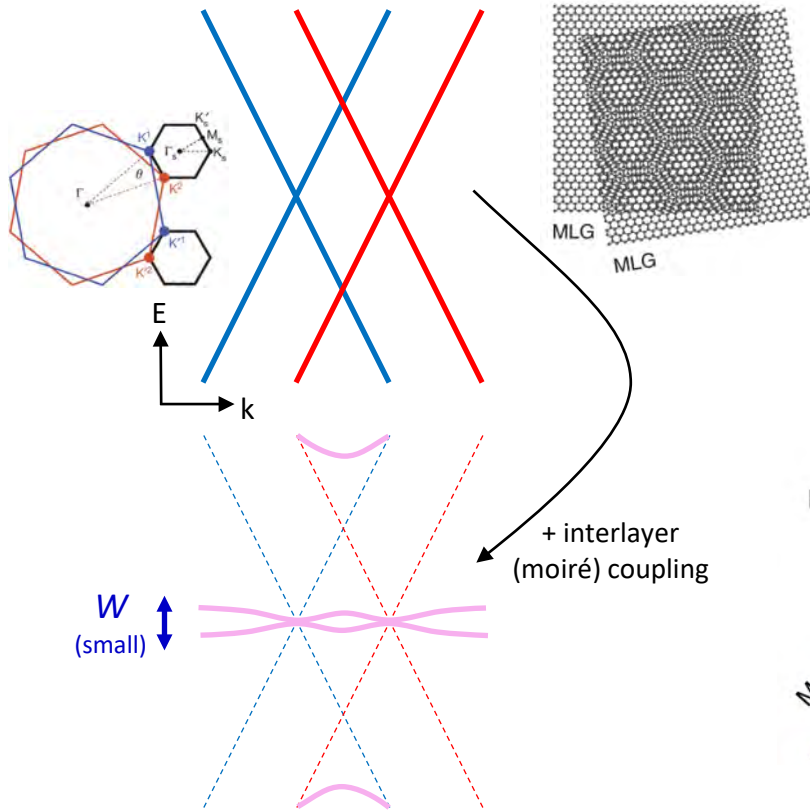
atomic stacking (and moiré)
→ *Berry-curved bands* and *topology*



Platforms for intermixed correlations and topology



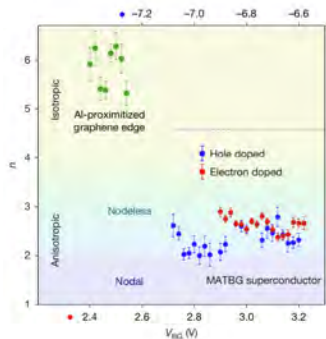
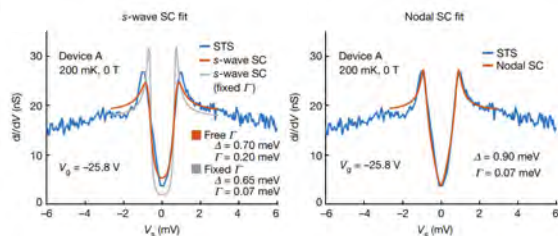
Summary of a “typical” MATBG phase diagram



Big open questions

superconductivity

What sets the pairing mechanism and gap symmetry?

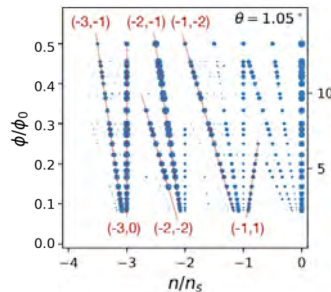
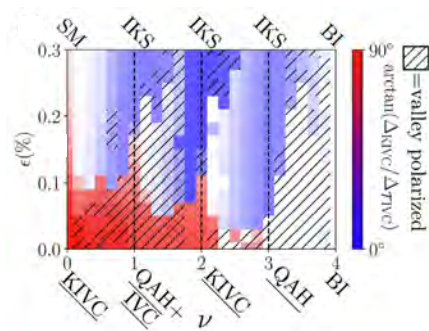


Evidence suggests unconventional superconductivity, but a definitive mechanism and pairing state remain unsettled.

Too big a question for today!

correlated ground states

What order and topology arise across the flat bands?

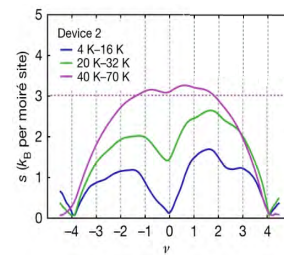
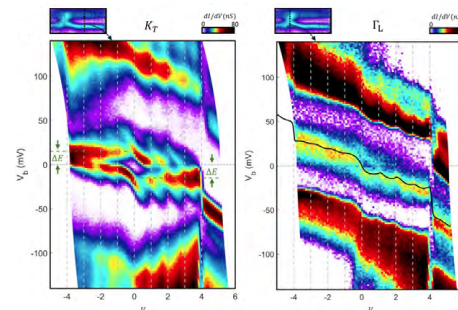


A broad picture is emerging, and many calculations and experiments agree on the main features.

Still, there are surprises.

normal state

What is the organizing principle of the high-temperature normal state (~100 K)?



We lack a fully quantitative theory for the high-temperature transport and thermodynamics.

There are suggestive links to heavy-fermion phenomenology with nontrivial topology, but we have more to learn.

Hartree-Fock predictions for MATBG ground states

PHYSICAL REVIEW LETTERS **128**, 156401 (2022)

PHYSICAL REVIEW X **14**, 021042 (2024)

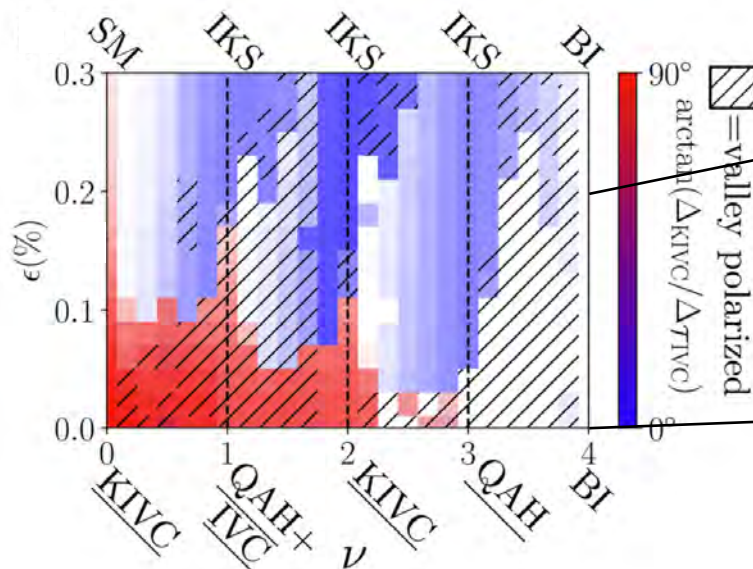
Global Phase Diagram of the Normal State of Twisted Bilayer Graphene

Glenn Wagner^{1,2}, Yves H. Kwan¹, Nick Bultinck^{1,3}, Steven H. Simon¹, and S. A. Parameswaran¹

¹Rudolf Peierls Centre for Theoretical Physics, Parks Road, Oxford OX1 3PU, United Kingdom

²Department of Physics, University of Zurich, Winterthurerstrasse 190, 8057 Zurich, Switzerland

³Department of Physics, Ghent University, Krijgslaan 281, 9000 Gent, Belgium

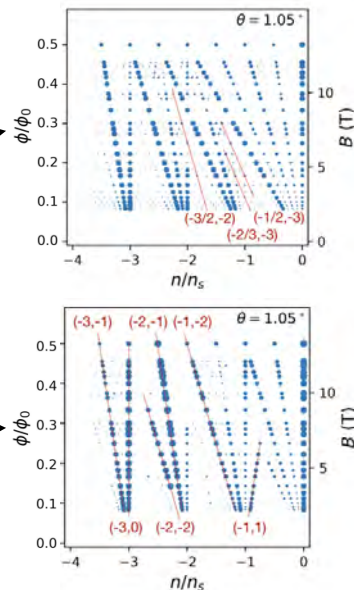


Theory of Correlated Chern Insulators in Twisted Bilayer Graphene

Xiaoyu Wang^{1,*} and Oskar Vafek^{1,2,†}

¹National High Magnetic Field Lab, Tallahassee, Florida 32310, USA

²Department of Physics, Florida State University, Tallahassee, Florida 32306, USA



- many closely competing ground states with varying spin, valley, and sublattice ordering
- symmetry-broken states evolve into a cascade of Chern insulators upon formation of Hofstadter subbands in a magnetic field
- ground state ordering is very sensitive to heterostrain between the two graphene sheets!

Chern insulator states tunable by doping

nature physics

Article

<https://doi.org/10.1038/s41567-025-02997-4>

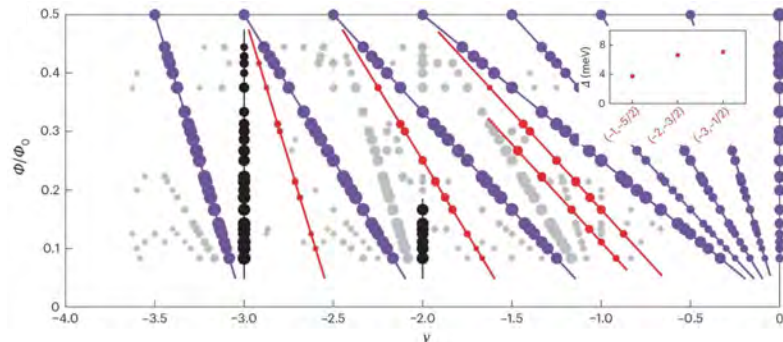
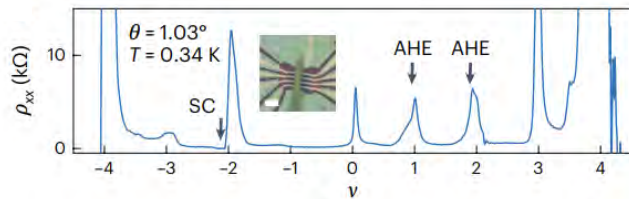
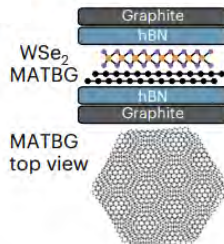
Strongly interacting Hofstadter states in magic-angle twisted bilayer graphene

Received: 1 August 2024

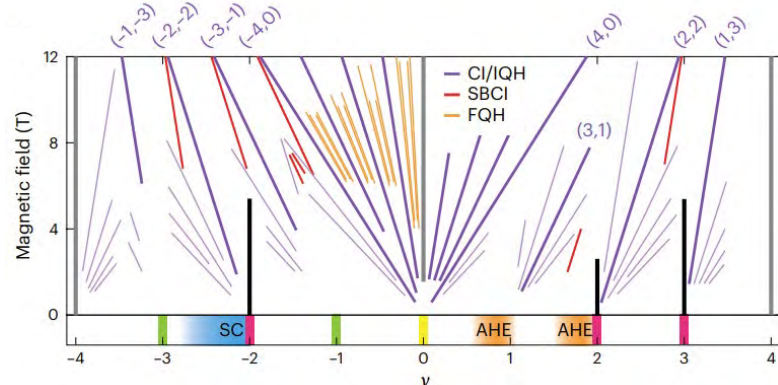
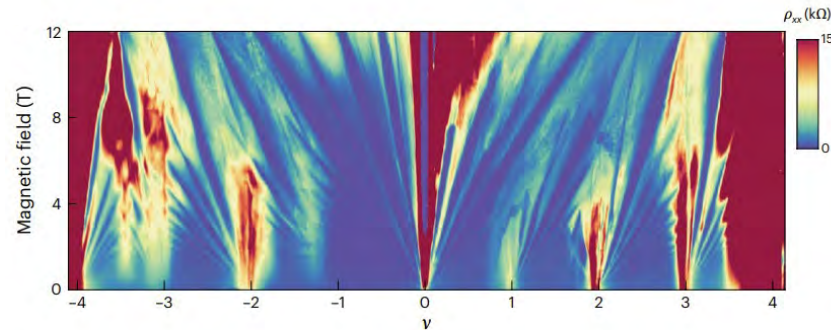
Accepted: 14 July 2025

Published online: 11 August 2025

Minhao He^{1,2,3}, Xiaoyu Wang^{1,2,3}, Jiaqi Cai¹, Jonah Herzog-Arbeitman⁴,
Ran Peng⁵, Takashi Taniguchi⁶, Kenji Watanabe⁶, Ady Stern⁷,
B. Andrei Bernevig⁸, Matthew Yankowitz^{1,7}, Oskar Vafeek^{1,8} &
Xiaodong Xu^{1,7}



- cascades of Chern states appear in a magnetic field
- some states break the moiré translational symmetry
- sequences broadly captured by Hartree-Fock calculations



Still a few surprises to understand

LETTERS

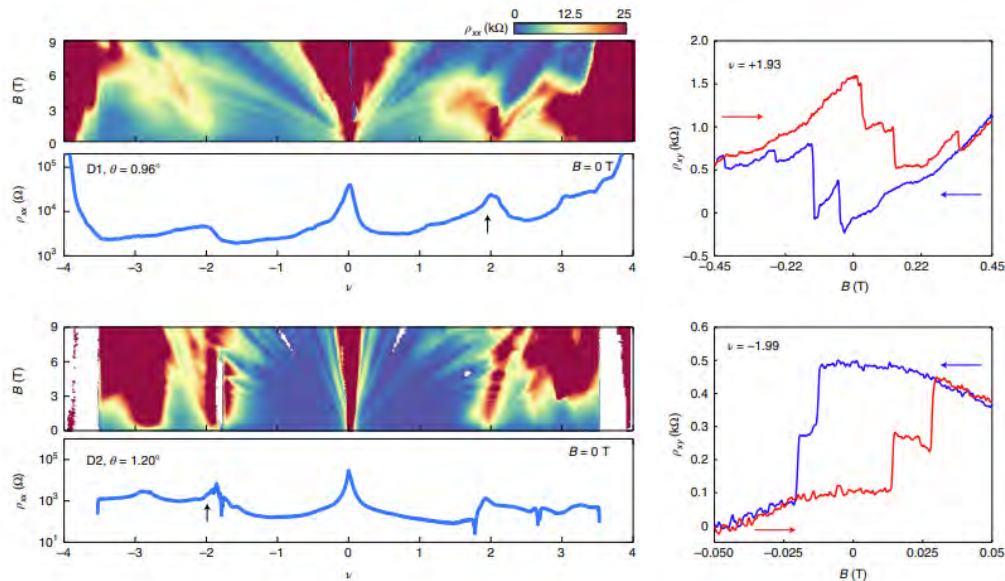
<https://doi.org/10.1038/s41567-022-01697-7>

nature
physics

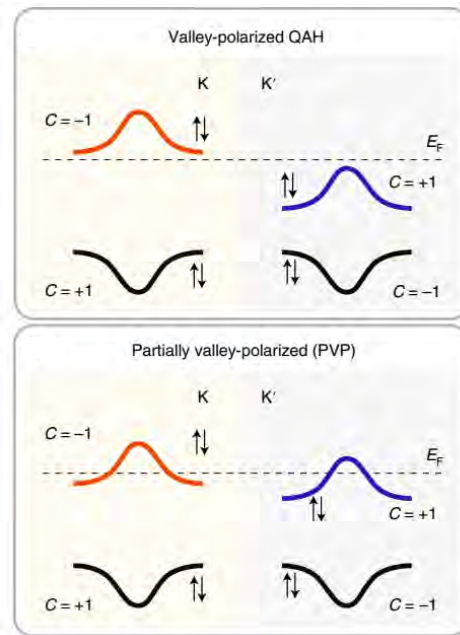
Check for updates

Anomalous Hall effect at half filling in twisted bilayer graphene

Chun-Chih Tseng^{1,5}, Xuetao Ma^{2,5}, Zhaoyu Liu^{1,5}, Kenji Watanabe³, Takashi Taniguchi⁴, Jiun-Haw Chu¹ and Matthew Yankowitz^{1,2}✉



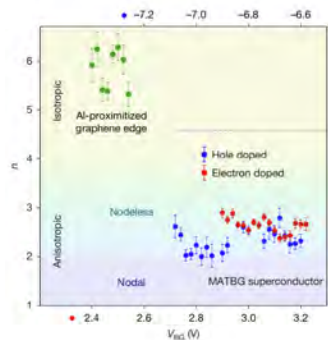
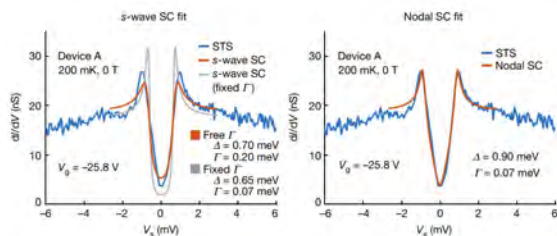
- intervalley Hund's coupling disfavors valley polarized states at $\nu = \pm 2$...
- ...so why do they appear here?



Big open questions

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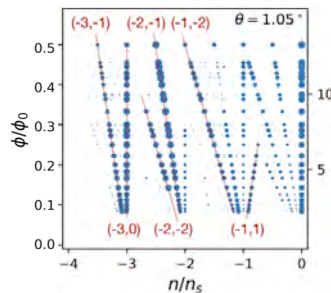
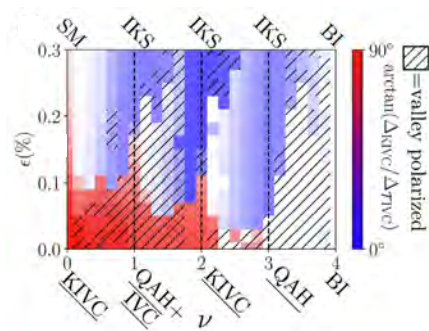


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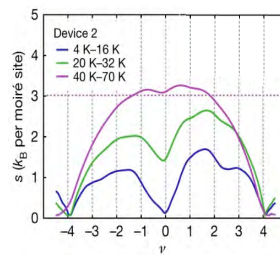
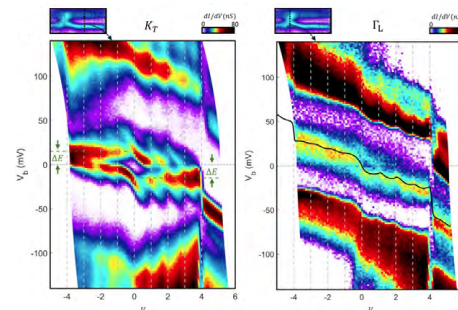


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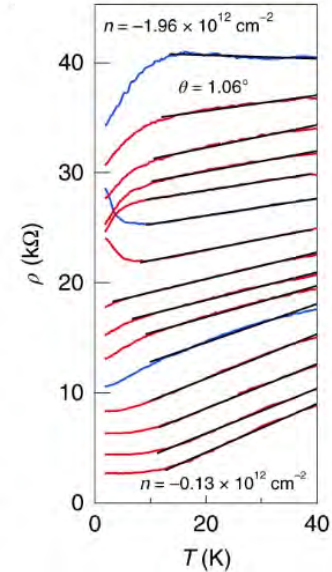


We lack a fully quantitative theory for the high-temperature transport and thermodynamics.

There are suggestive links to heavy-fermion phenomenology with nontrivial topology, but we have more to learn.

Unusual transport in the normal state of MATBG

H. Polshyn, M. Yankowitz et al.,
Nature Physics 15, 1011 (2019)



See also: Y. Cao et al., PRL 124, 076801 (2020)
A. Jaoui et al., Nature Physics 18, 633 (2022)

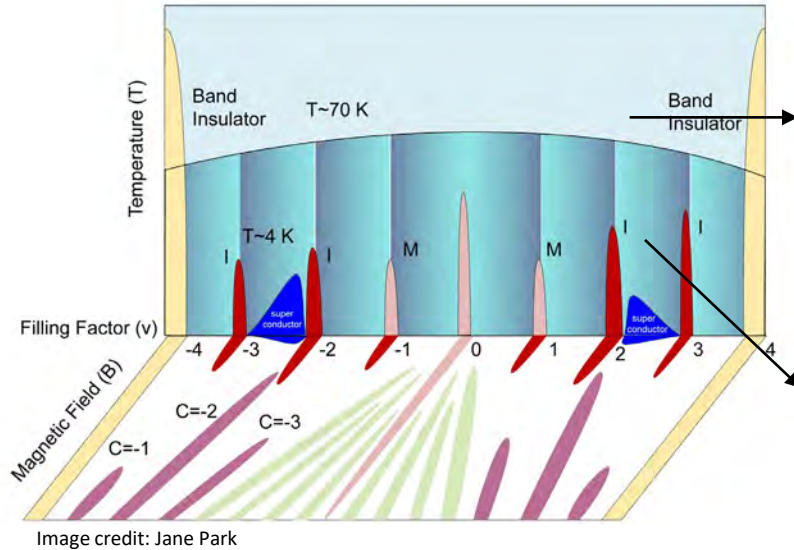
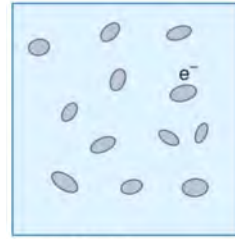
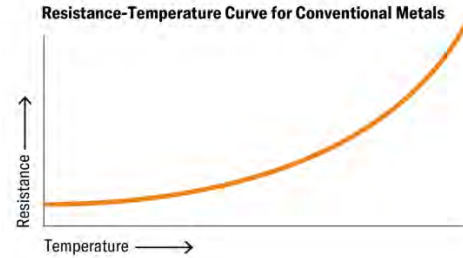
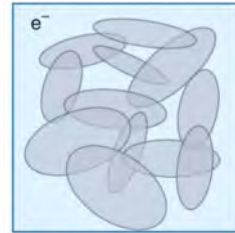
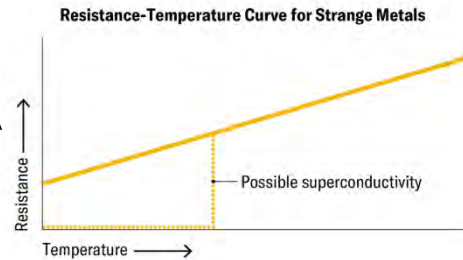


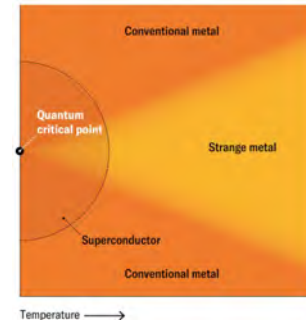
Image credit: Jane Park



Normal Metal



Strange Metal



T. Stauber et al., Nature
Physics 18, 619 (2022)

D. Natelson, Scientific
American 330, 42 (2024)

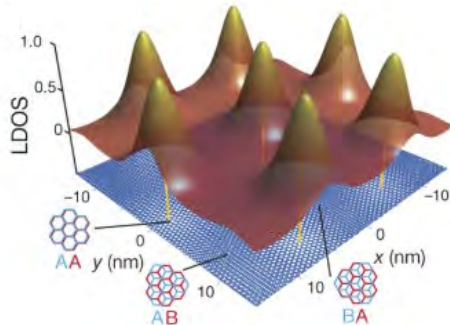
related strange metals and superconductivity in:

cuprates, iron-based superconductors, nickelates, ruthenates, Kagome metals, heavy-fermion Kondo lattices, organic charge-transfer salts, etc.

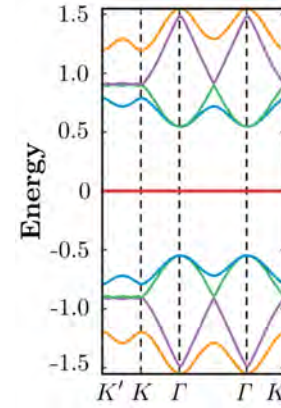
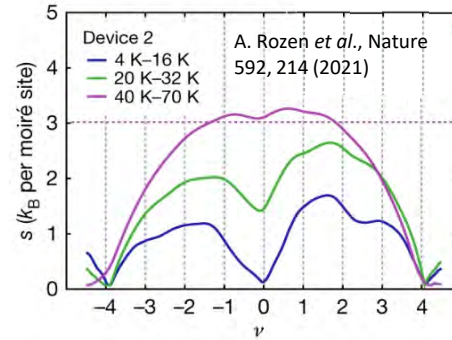
Local or nonlocal physics in MATBG?

local description:

Mott insulators on a triangular lattice of AA-stacked sites

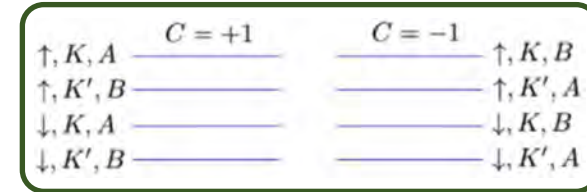


Y. Cao *et al.*, Nature 556, 43 (2018)



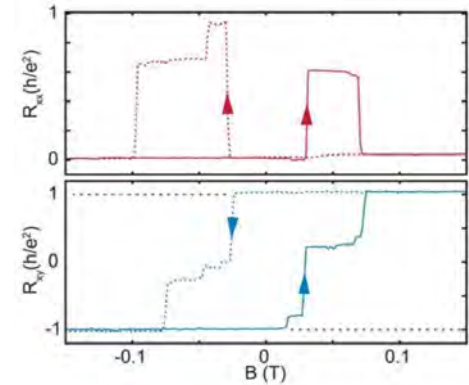
nonlocal description:
analog of quantum Hall ferromagnets

G. Tarnopolsky *et al.*, PRL 122, 106405 (2019)



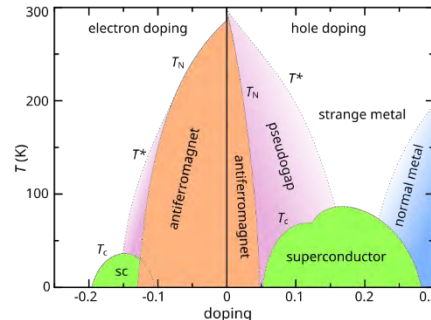
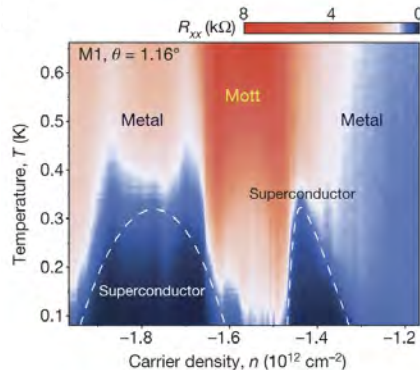
Serlin *et al.*, Science 367, 900 (2020)

see also: Sharpe *et al.*, Science 365, 605 (2019)



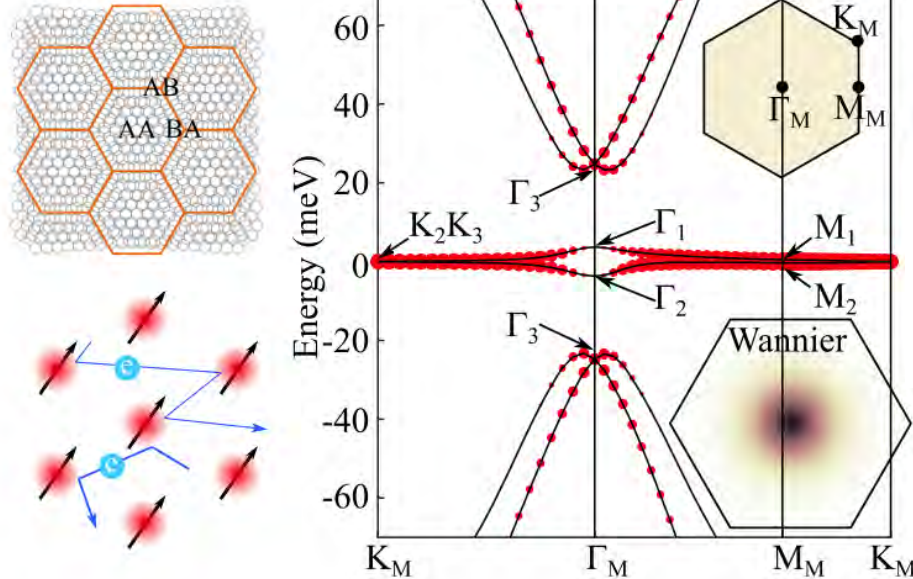
reminiscent of
superconductivity from
doped Mott insulators

topological states require
delocalized electrons



Topological heavy fermion model

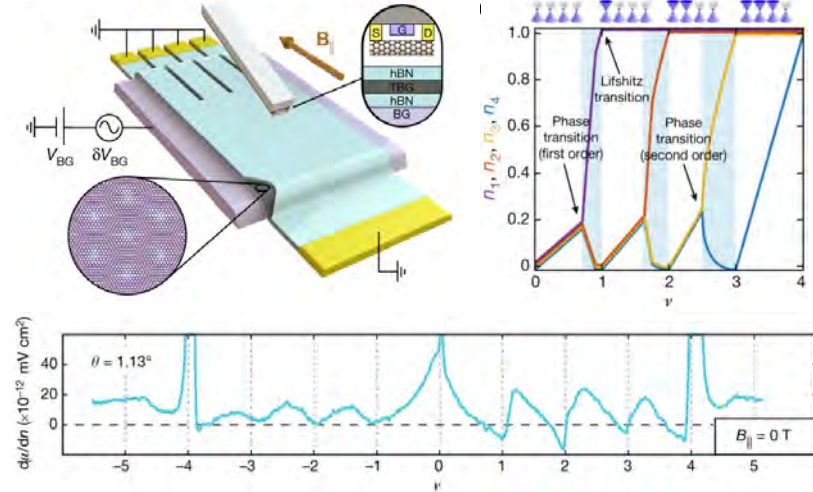
Z.-D. Song and B. A. Bernevig, PRL 129, 047601 (2022)



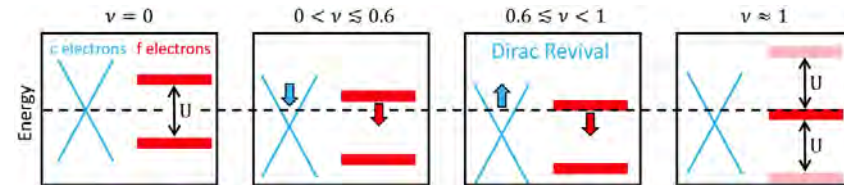
- topology/symmetry of the moiré lattice obstructs the formation of exponentially localized Wannier orbitals
- construct bands by hybridizing localized moments on the AA-stacking sites with dispersive bands from AB/BA-sites

**low-temperature:
symmetry breaking**

U. Zondiner *et al.*, Nature 582, 203 (2020)



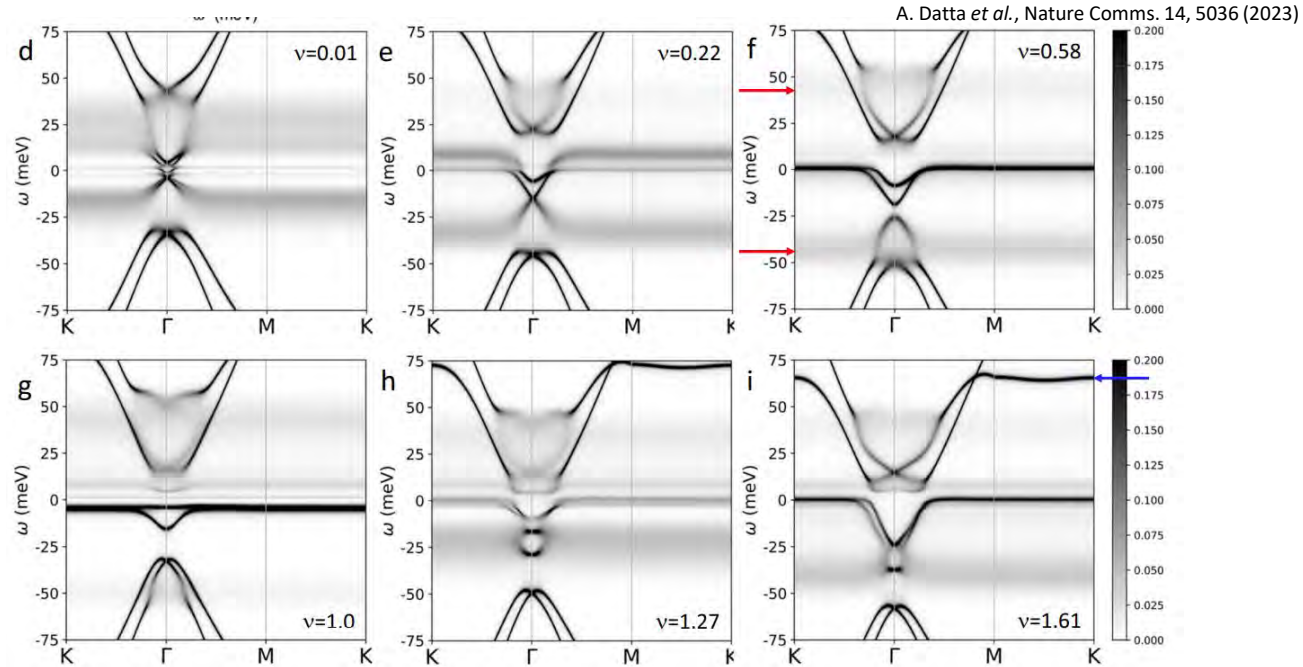
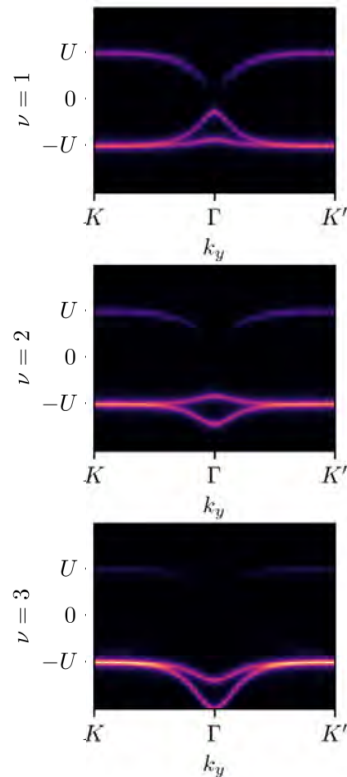
high temperature: heavy fermion physics



J. Xiao *et al.*, arXiv:2506.20738 (2025)

No matter the exact model, do many-body calculations

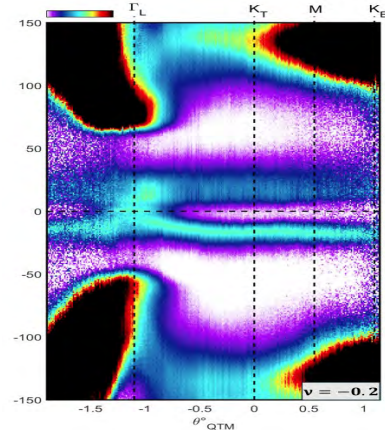
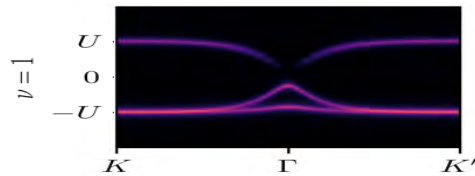
Hartree-Fock + dynamical mean-field theory (DMFT) \rightarrow spectral function



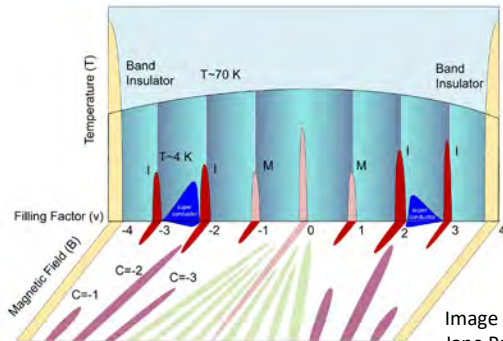
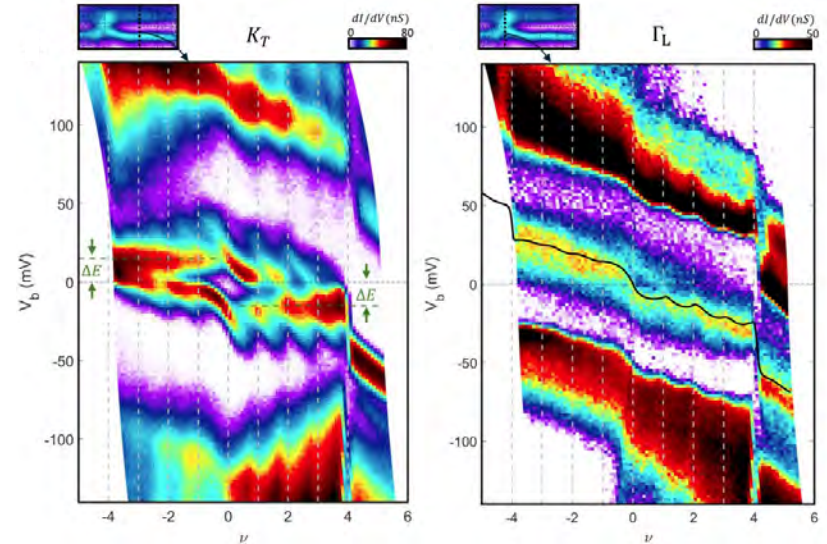
calculations generally agree on flat, spectrally-incoherent “Hubbard bands” broadened by $\sim 2U$ almost everywhere across the moiré Brillouin zone, along with dispersive, sharply defined quasiparticles near Γ

Key recent findings from tunneling measurements

interactions profoundly reshape the bands at the magic angle



quantum twisting microscope measurements may help to unite the localized and extended electronic behavior with a topological heavy fermion framework



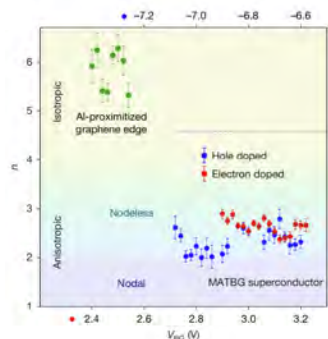
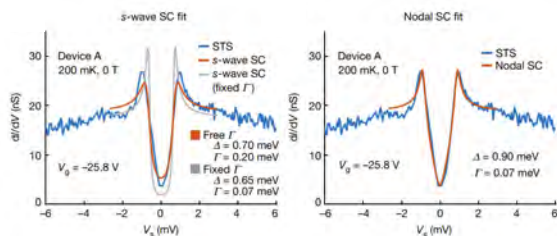
what happens < 4 K once isospin symmetries break, and correlated and superconducting gaps open?

Image credit:
Jane Park

Big open questions

superconductivity

What sets the pairing mechanism and gap symmetry?

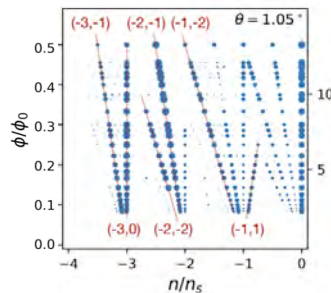
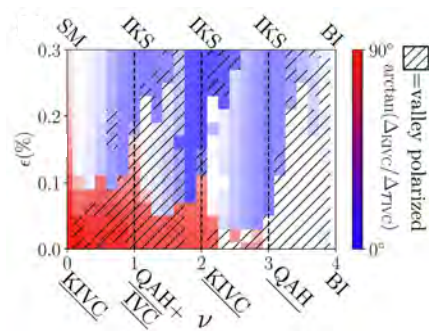


Evidence suggests unconventional superconductivity, but a definitive mechanism and pairing state remain unsettled.

Too big a question for today!

correlated ground states

What order and topology arise across the flat bands?

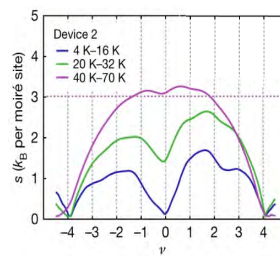
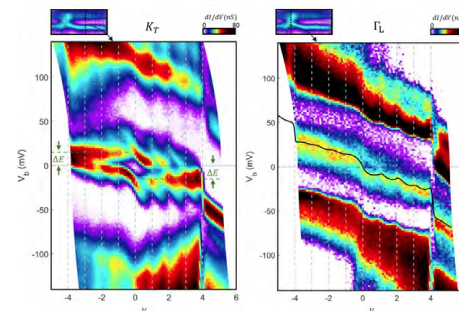


A broad picture is emerging, and many calculations and experiments agree on the main features.

Still, there are surprises.

normal state

What is the organizing principle of the high-temperature normal state (~ 100 K)?

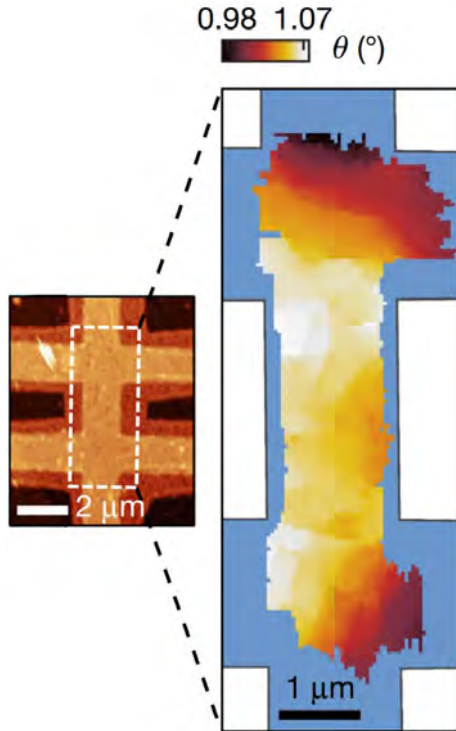


We lack a fully quantitative theory for the high-temperature transport and thermodynamics.

There are suggestive links to heavy-fermion phenomenology with nontrivial topology, but we have more to learn.

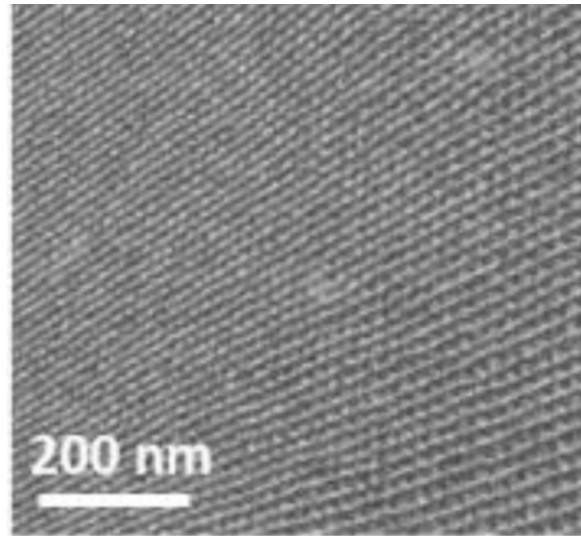
A challenge: prevalence of unintended strain

scanning SQUID



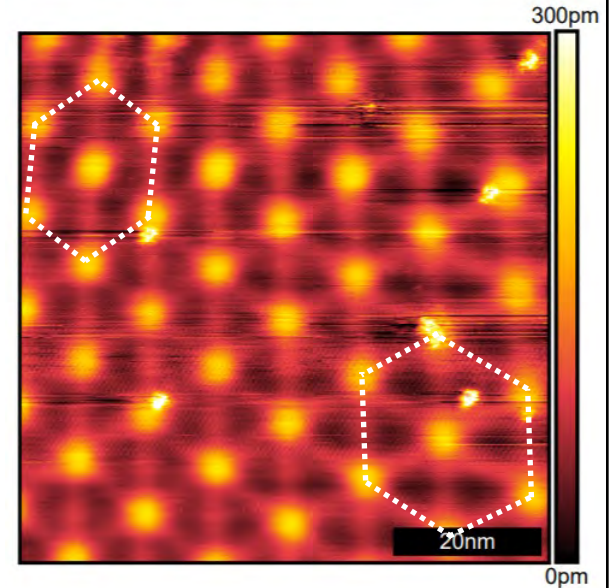
Zeldov group, Nature (2020)

TEM



Kim group, Nature Materials (2019)

STM



Pasupathy group, Nature (2019)

Unintended strain complicates device behavior

Review

Reproducibility in the fabrication and physics of moiré materials

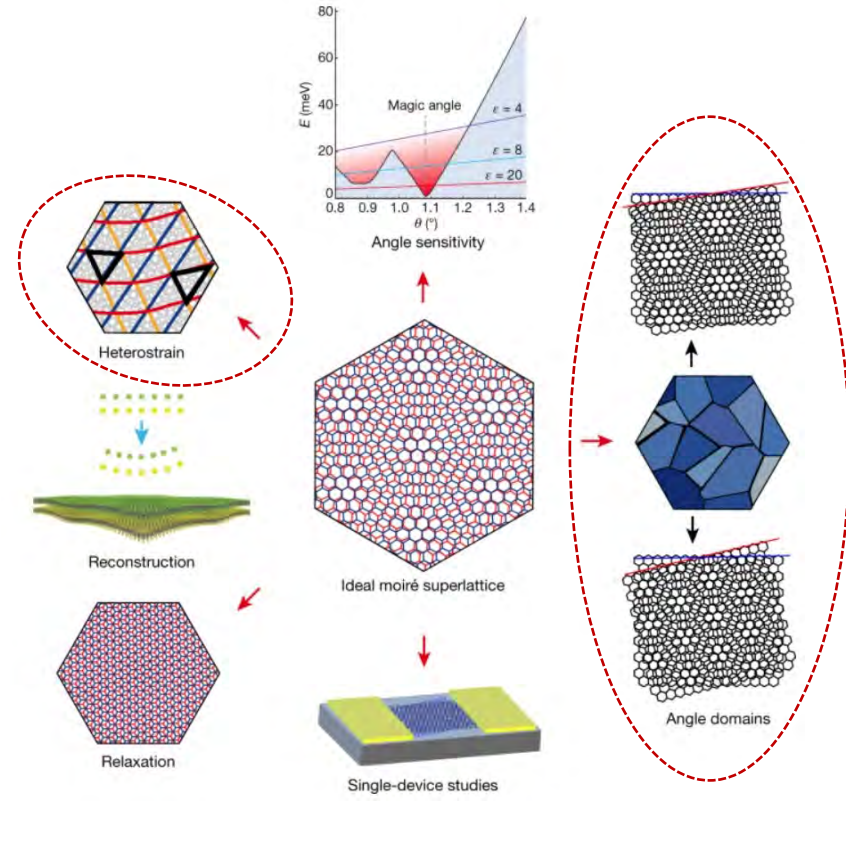
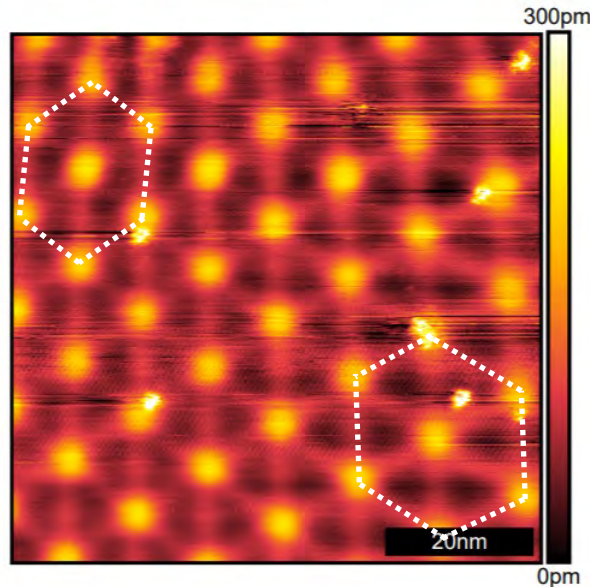
<https://doi.org/10.1038/s41586-021-04173-z>

Received: 16 March 2021

Accepted: 21 October 2021

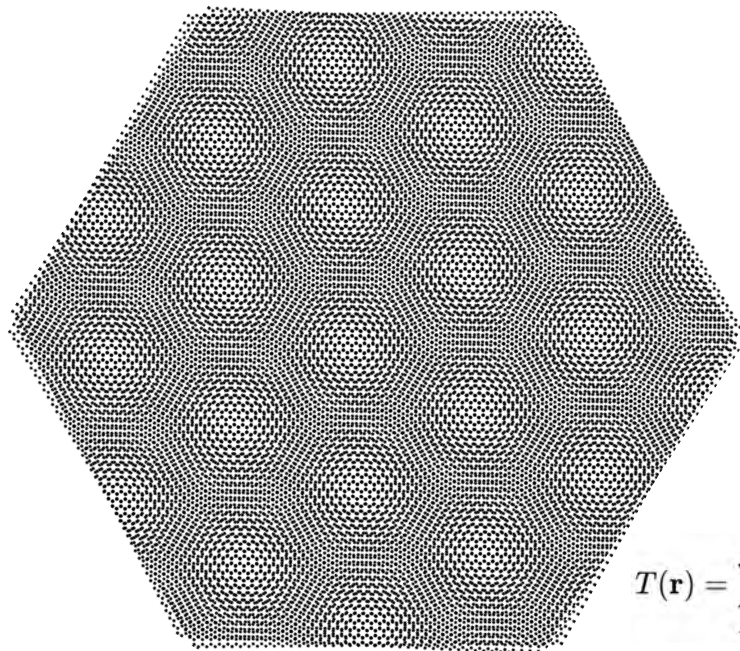
Published online: 2 February 2022

Chun Ning Lau^{1,2}, Marc W. Bockrath¹, Kin Fai Mak² & Fan Zhang³

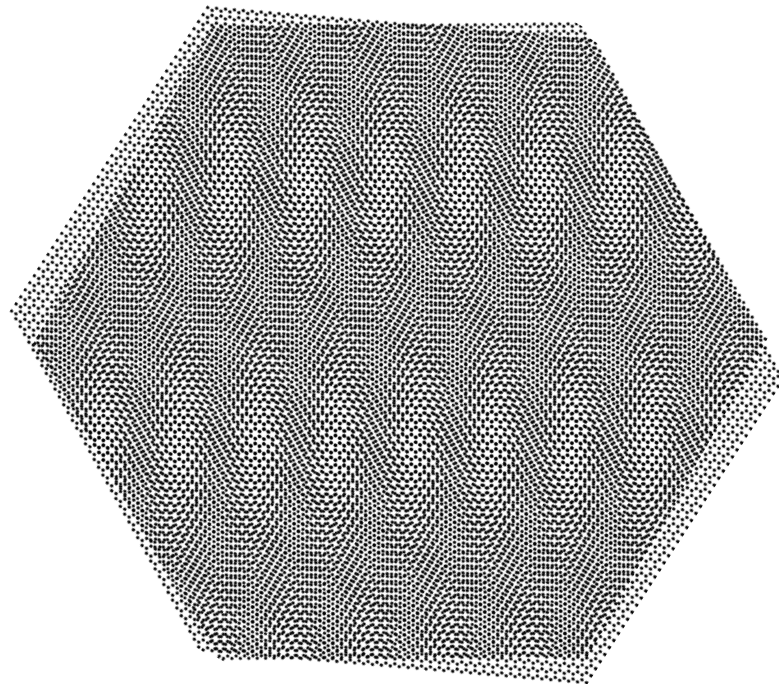


Strain modifies the moiré Hamiltonian

unstrained



heterostrain

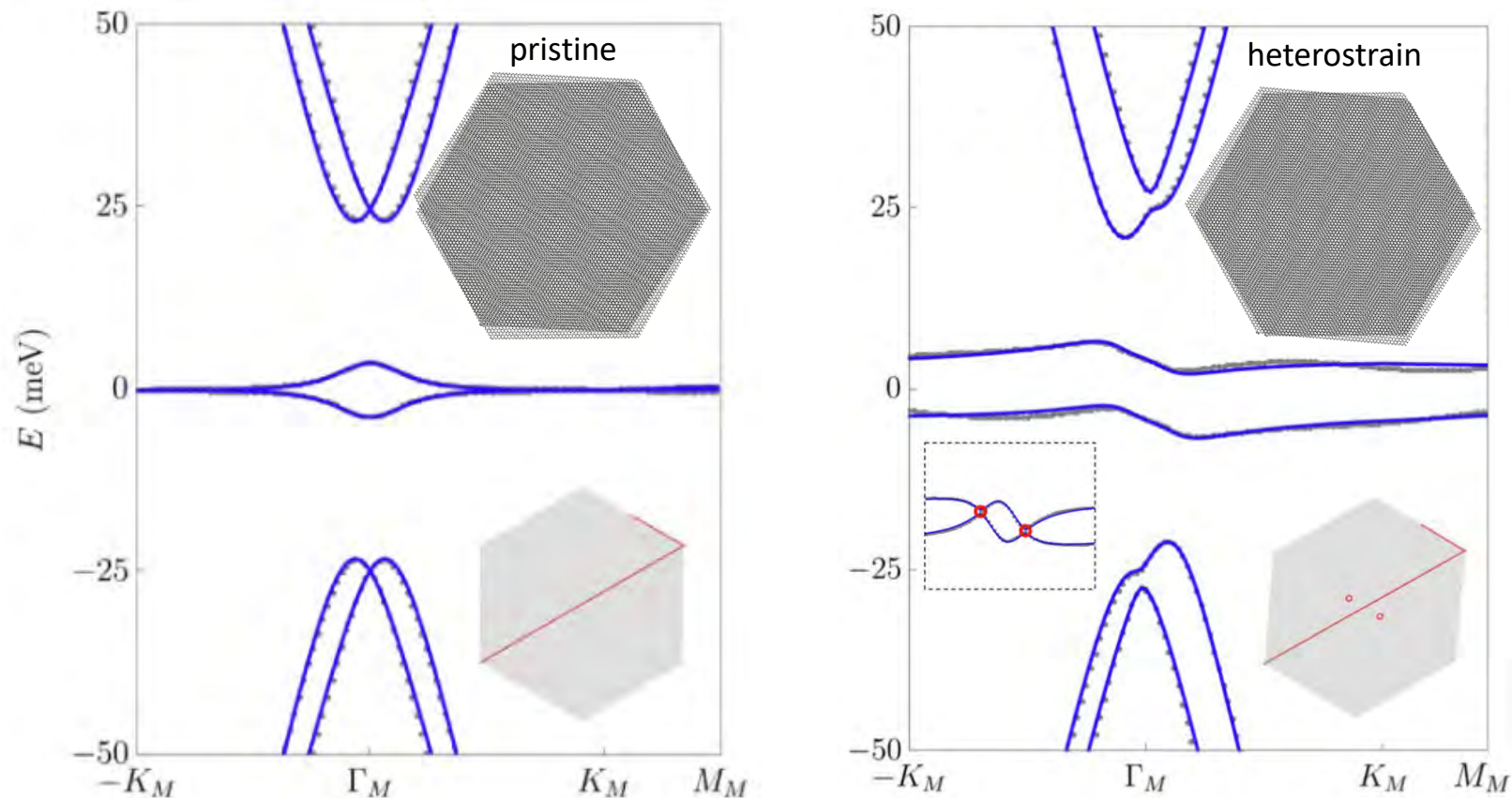


$$T(\mathbf{r}) = \sum_{j=1}^3 T_j e^{i\mathbf{q}_j \cdot \mathbf{r}}$$

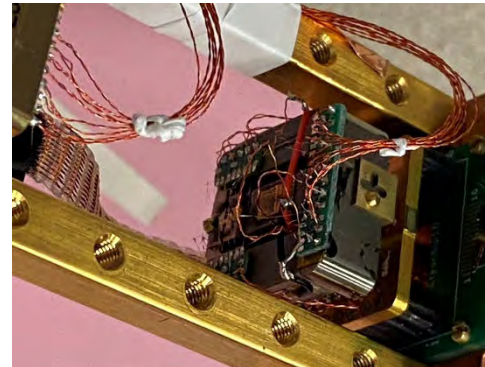
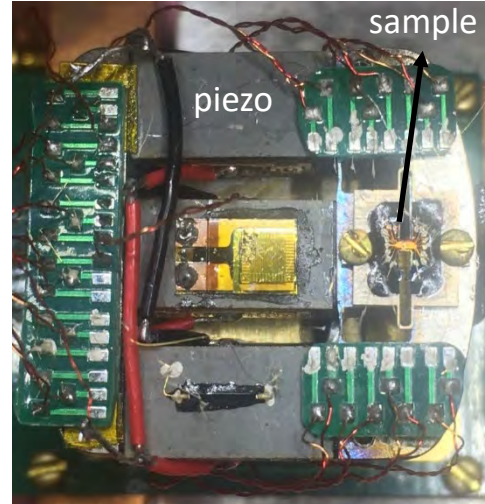
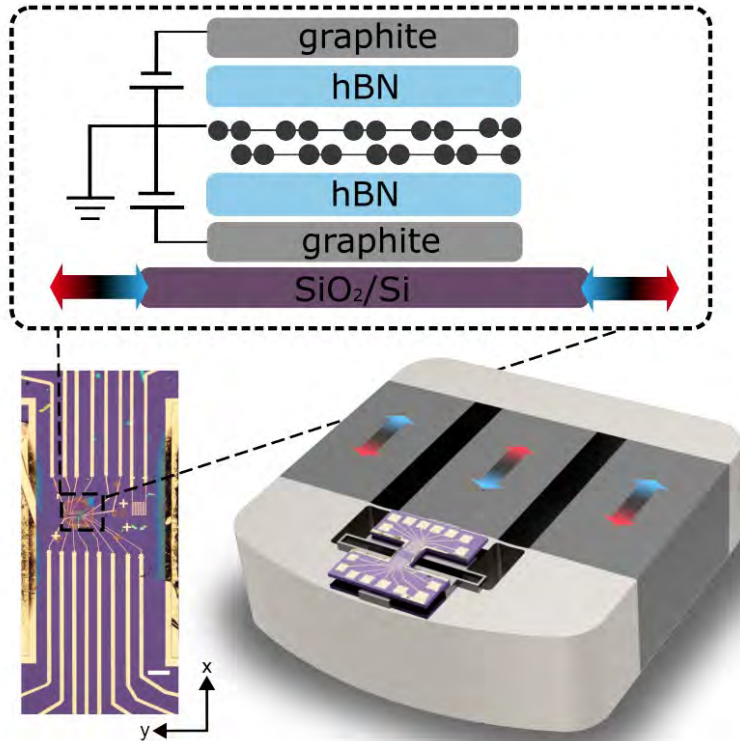
$$h_{\text{BM}}^K(\mathbf{r}) = \begin{pmatrix} -iv_F \nabla \cdot \boldsymbol{\sigma} & T(\mathbf{r}) \\ T^\dagger(\mathbf{r}) & -iv_F \nabla \cdot \boldsymbol{\sigma} \end{pmatrix}$$

$$\mathbf{q}_j \rightarrow \mathbf{q}_j(\varepsilon) = \mathbf{q}_j^{(0)} + \delta \mathbf{q}_j(\varepsilon)$$

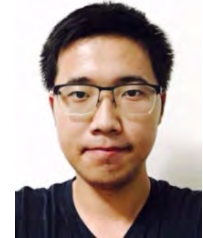
Effect of heterostrain on the moiré bands



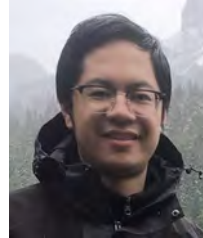
Directly controlling strain in vdW devices



Dr. Xuetao Ma



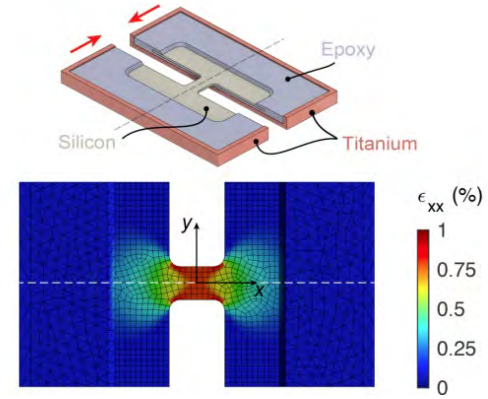
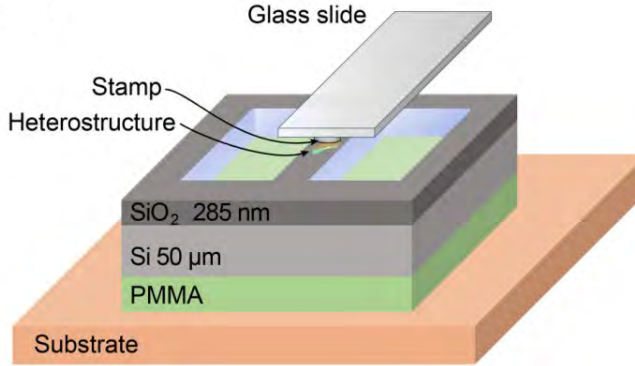
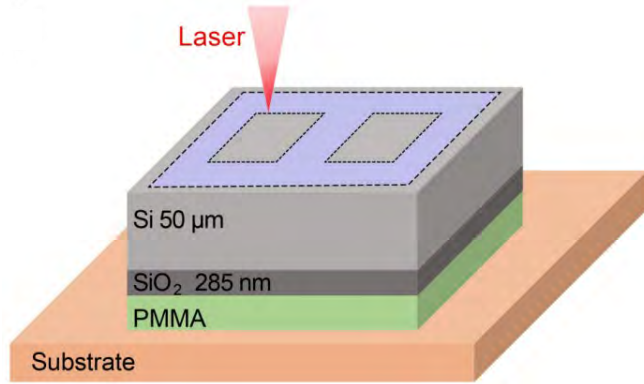
Dr. Zhaoyu Liu



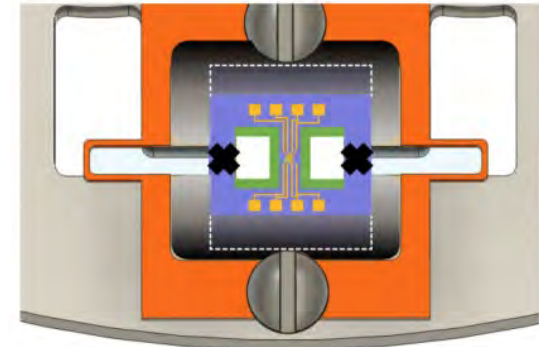
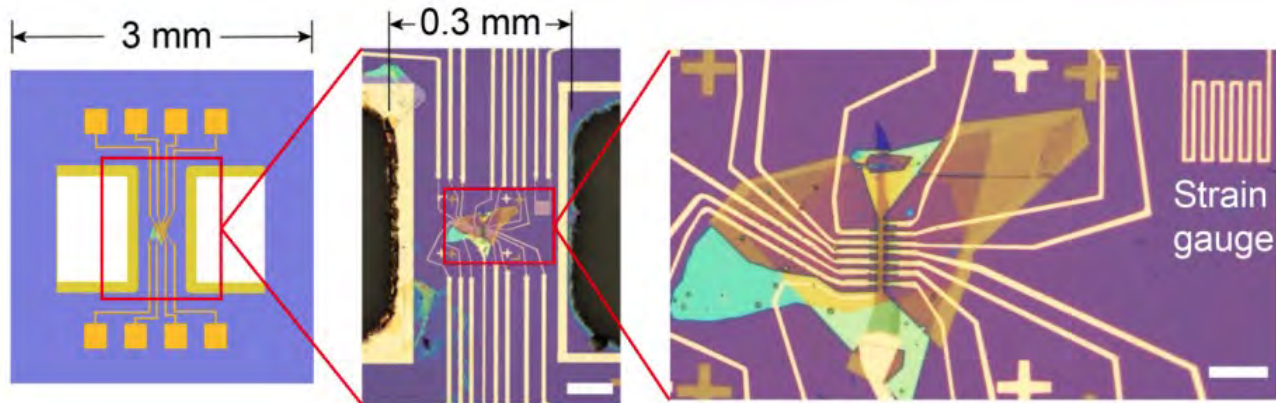
Prof. Jiun-Haw Chu



Achieving uniaxial strain in complex vdW devices

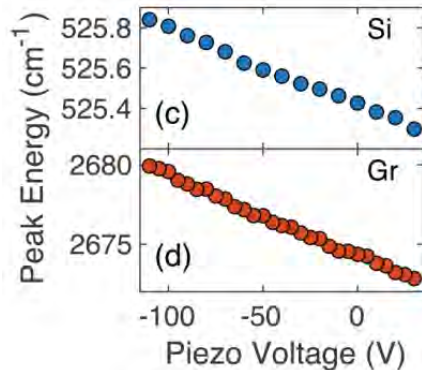
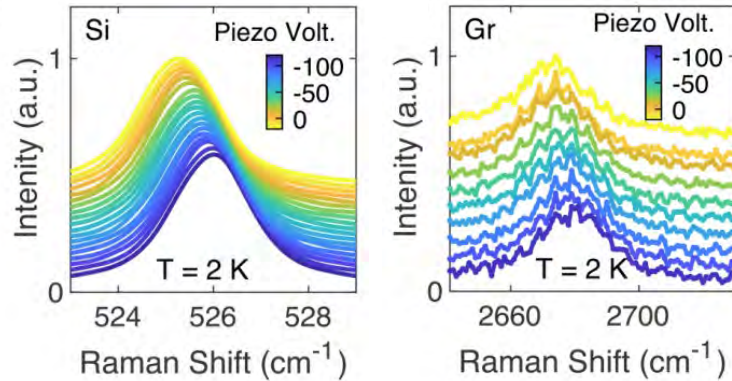


Z. Liu, M. Yankowitz, J.-H. Chu *et al.*, J. Appl. Phys. 135, 204306 (2024)

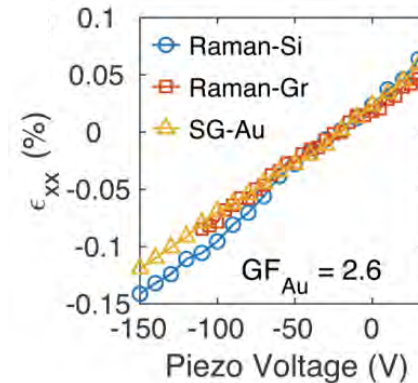
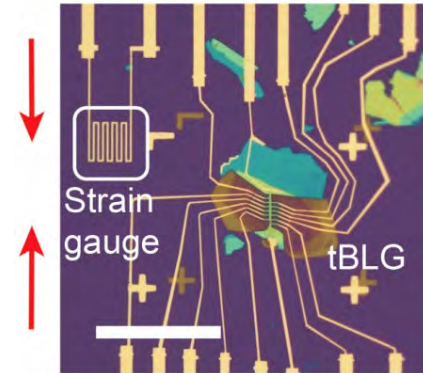


Measuring the induced strain

optically:



electrically:



Giant elastoresistance in magic-angle twisted graphene

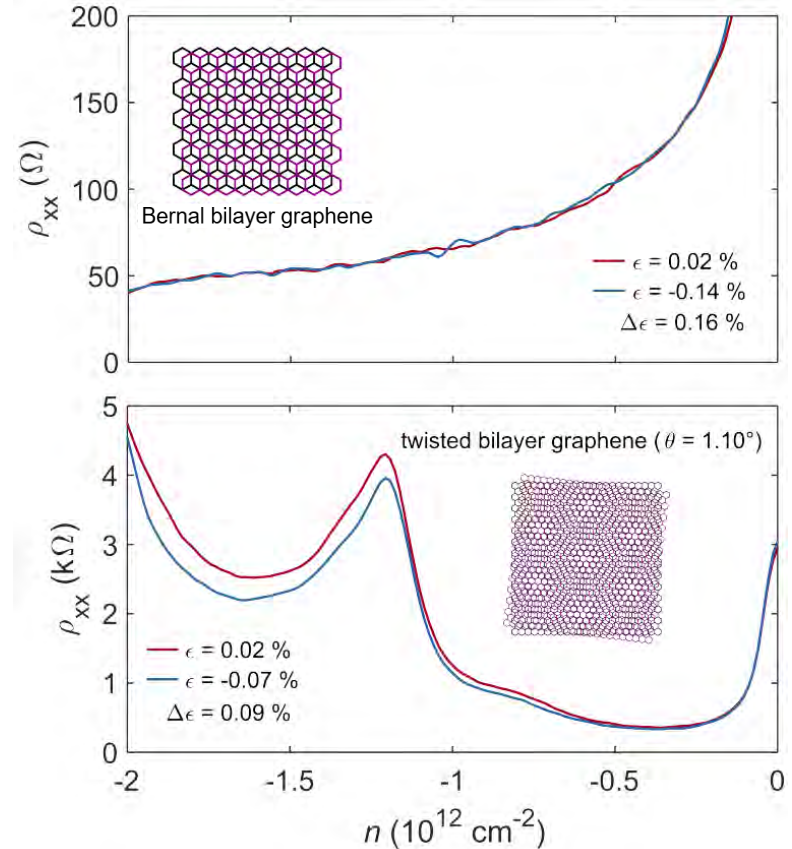
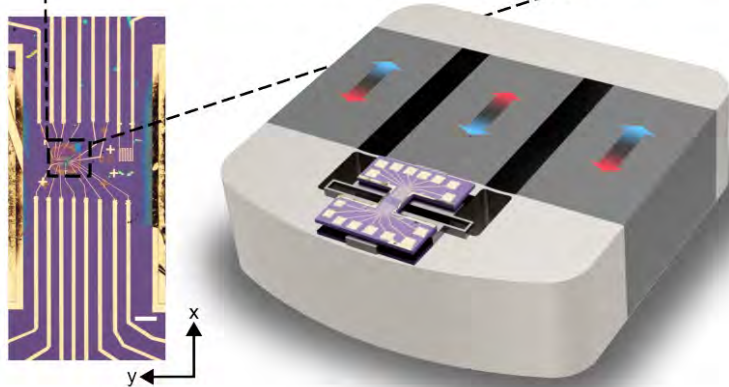
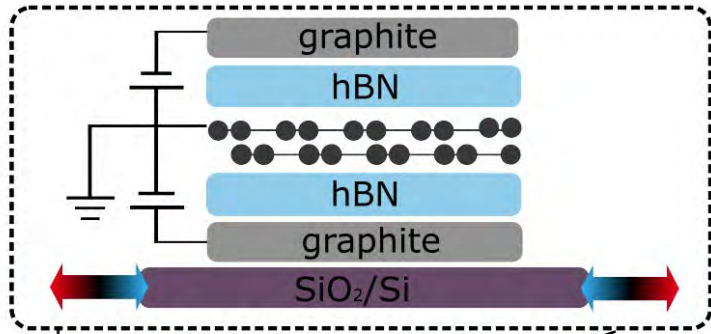


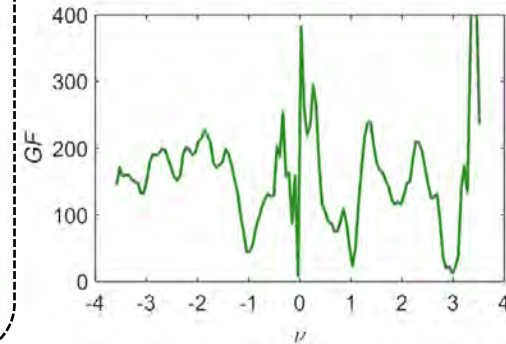
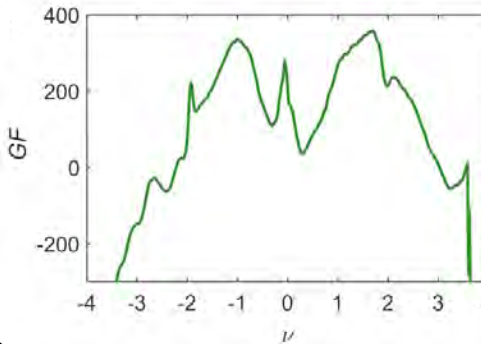
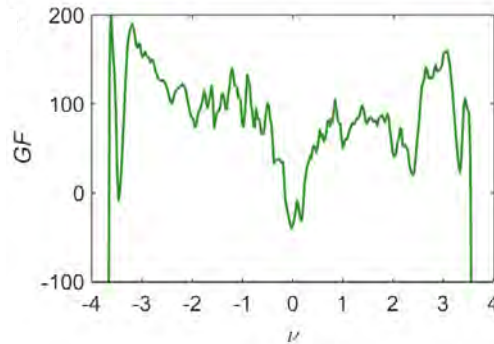
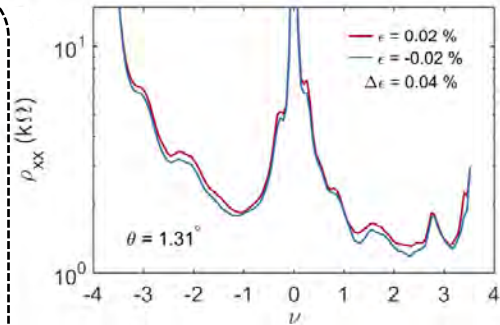
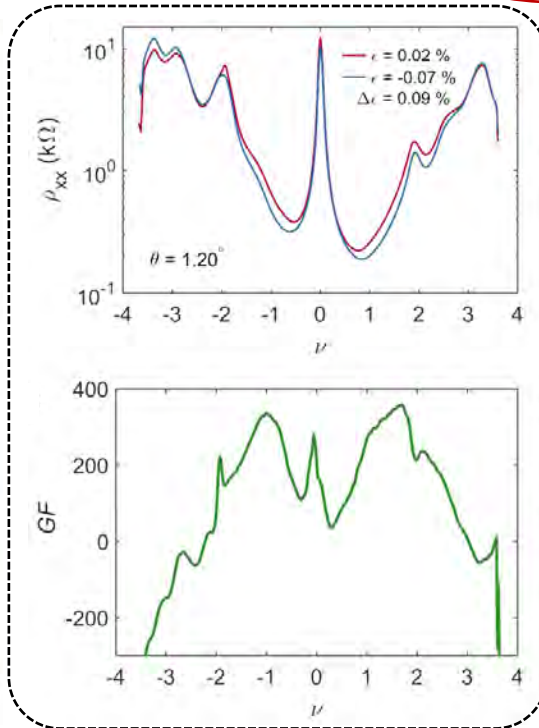
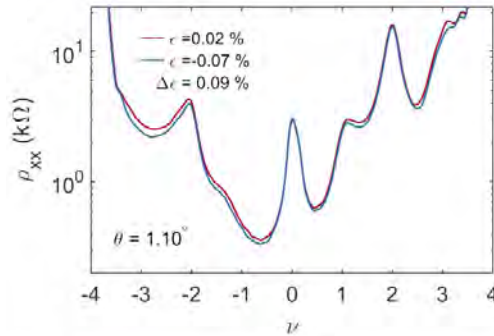
Figure of merit: gauge factor

$$GF = \frac{\Delta R/R}{\Delta L/L} = \frac{\Delta R/R}{\Delta \epsilon}$$

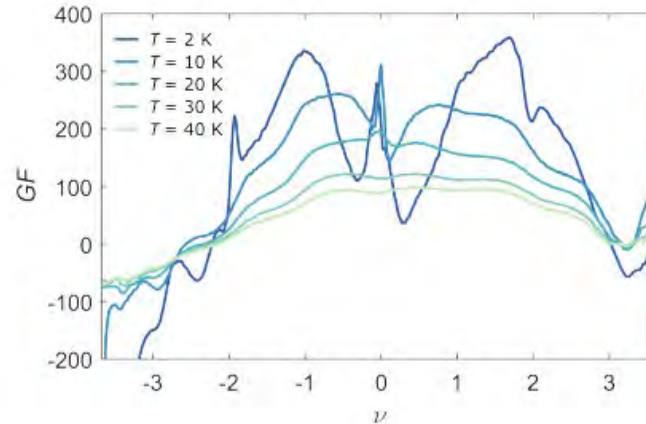
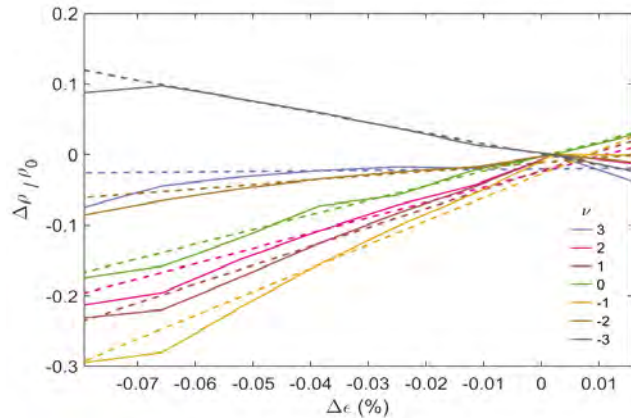
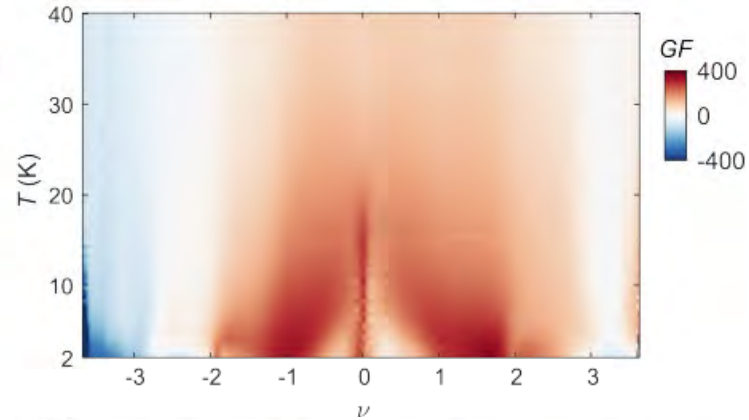
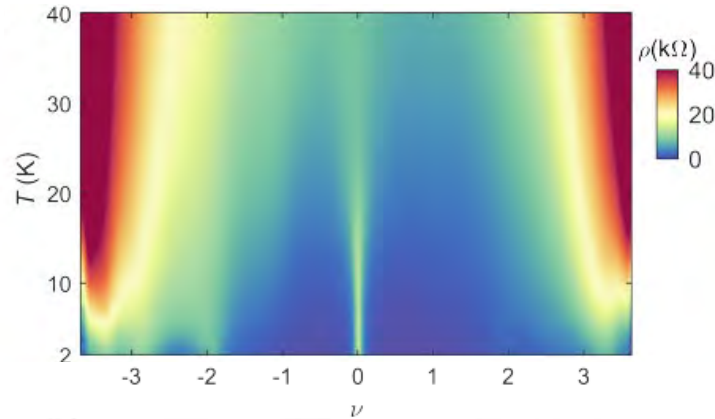
$$GF = 1 + 2\nu_P + \frac{\Delta \rho / \rho}{\epsilon}$$

electronic

geometric



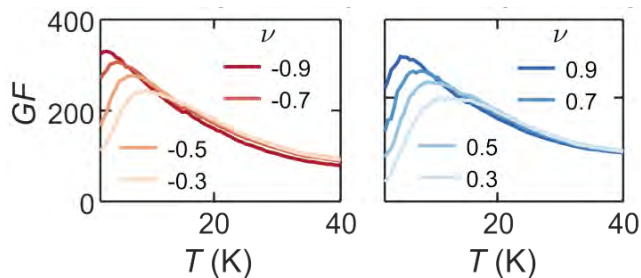
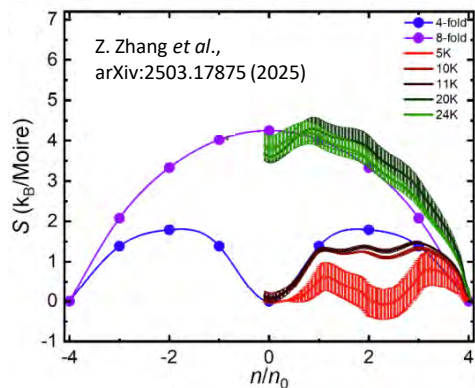
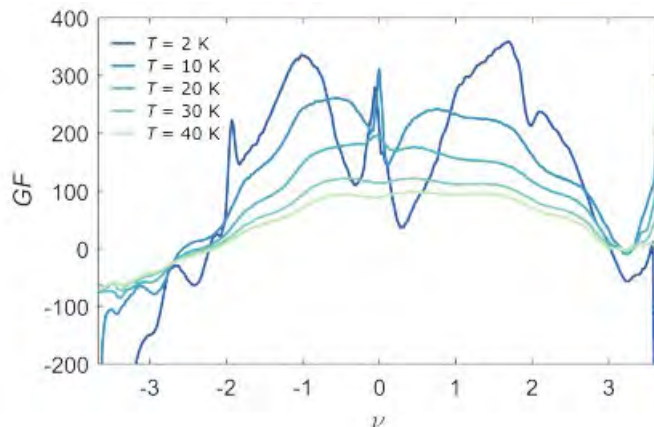
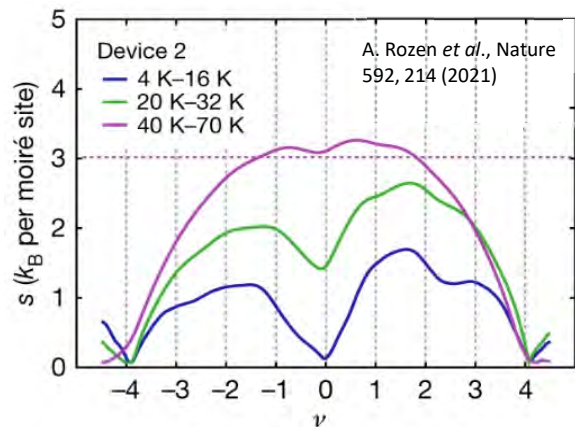
Temperature dependence of the elastoresistance



Key observations:

- giant gauge factor
- modulated by doping
- large temperature dependence

Large electronic entropy in MATBG



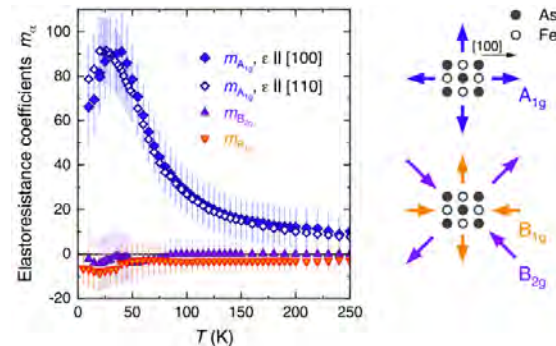
weakly diverging elastoresistance near the charge neutrality point

Dominant In-Plane Symmetric Elastoresistance in CsFe_2As_2

P. Wiecki¹, A.-A. Haghighirad¹, F. Weber¹, M. Merz¹, R. Heid¹, and A. E. Böhmer¹
Karlsruhe Institute of Technology, Institute for Quantum Materials and Technologies, 76021 Karlsruhe, Germany

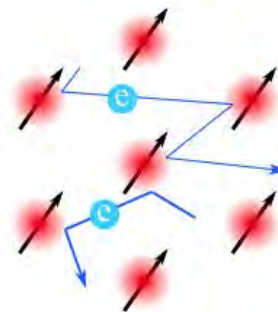
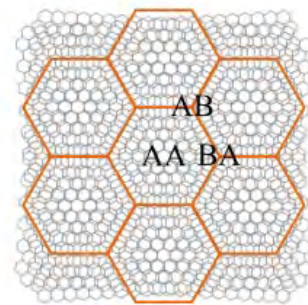
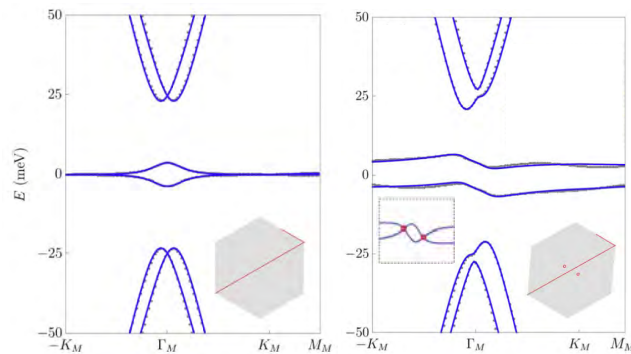
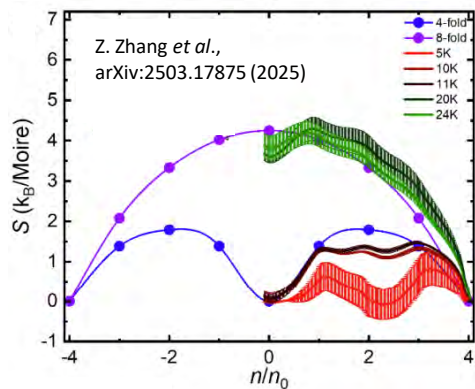
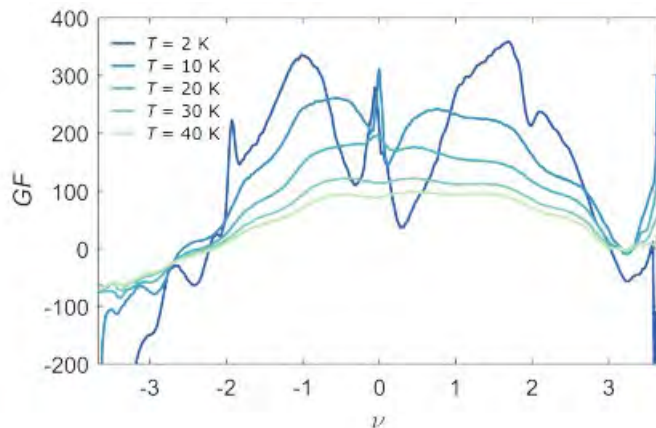
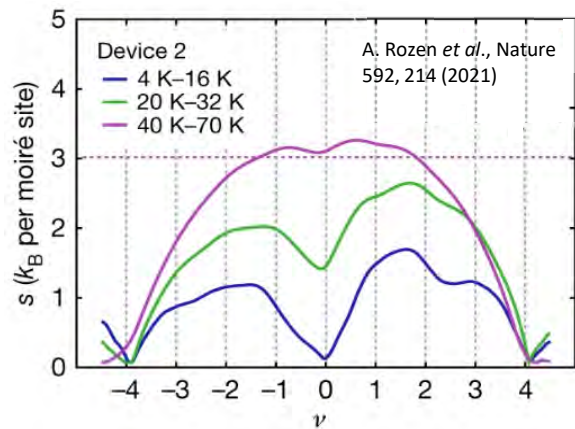
(Received 20 May 2020; revised 15 July 2020; accepted 24 September 2020; published 28 October 2020)

We study the elastoresistance of the highly correlated material CsFe_2As_2 in all symmetry channels. Neutralizing its thermal expansion by means of a piezoelectric-based strain cell is demonstrated to be essential. The elastoresistance response in the in-plane symmetric channel is found to be large, while the response in the symmetry-breaking channels is weaker and provides no evidence for a divergent nematic susceptibility. Rather, our results can be interpreted naturally within the framework of a coherence-incoherence crossover, where the low-temperature coherent state is sensitively tuned by the in-plane atomic distances.



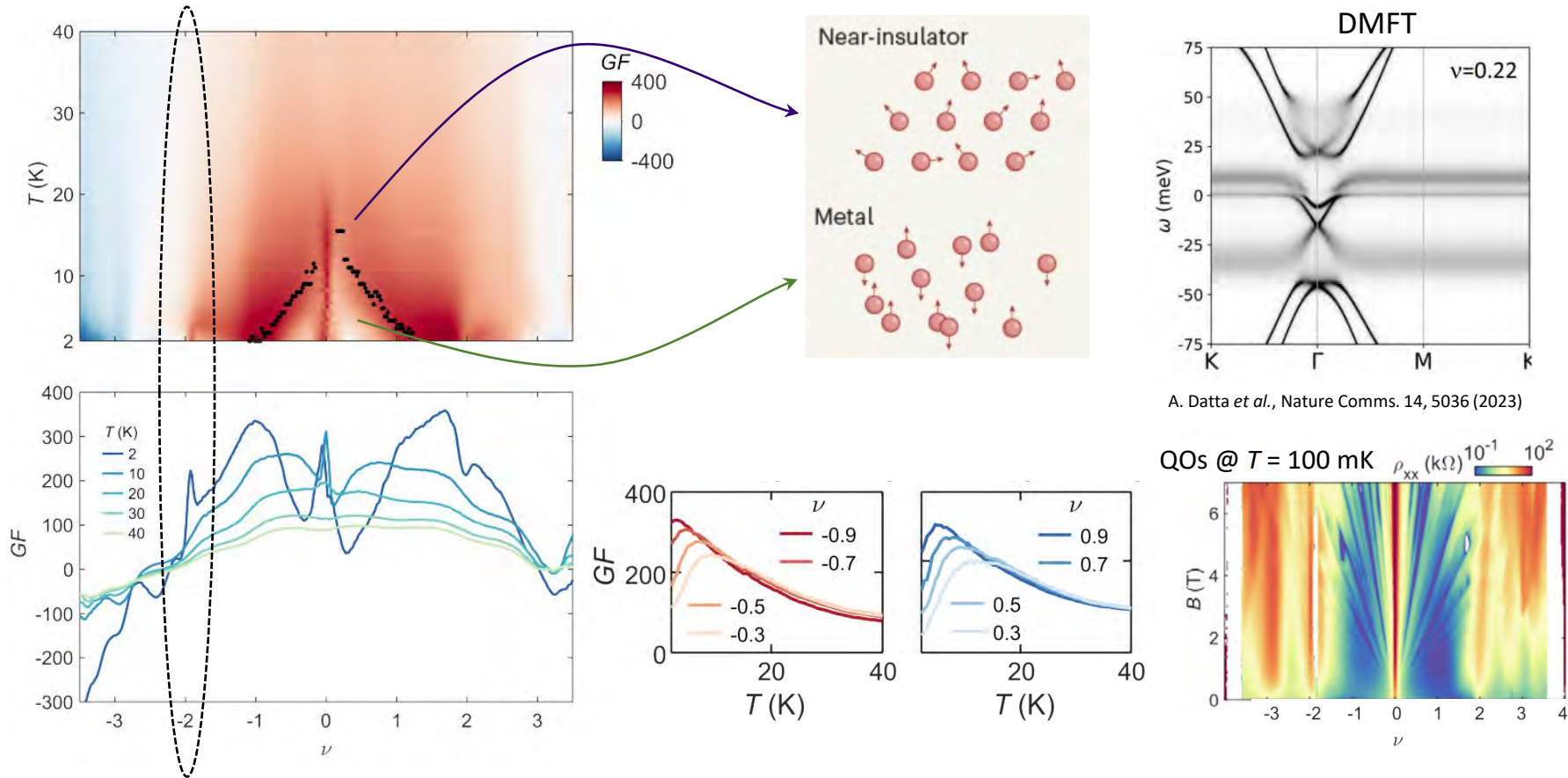
related observations in iron-based superconductors with weak nematicity

Large electronic entropy in MATBG



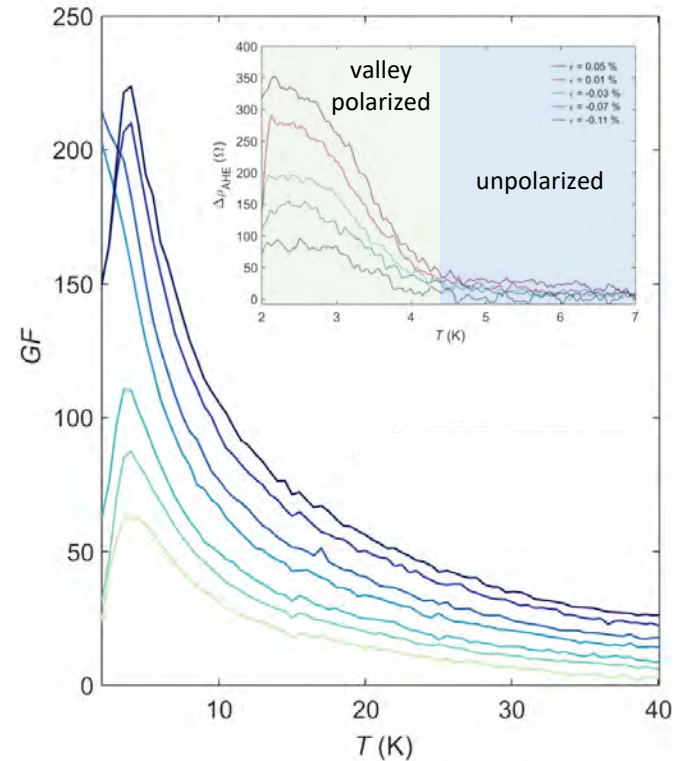
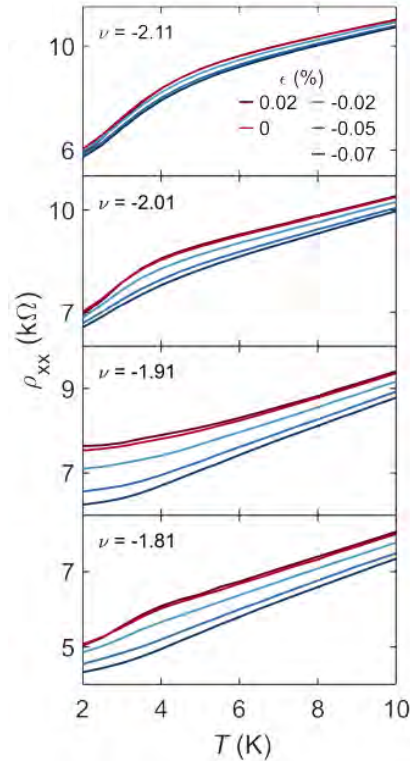
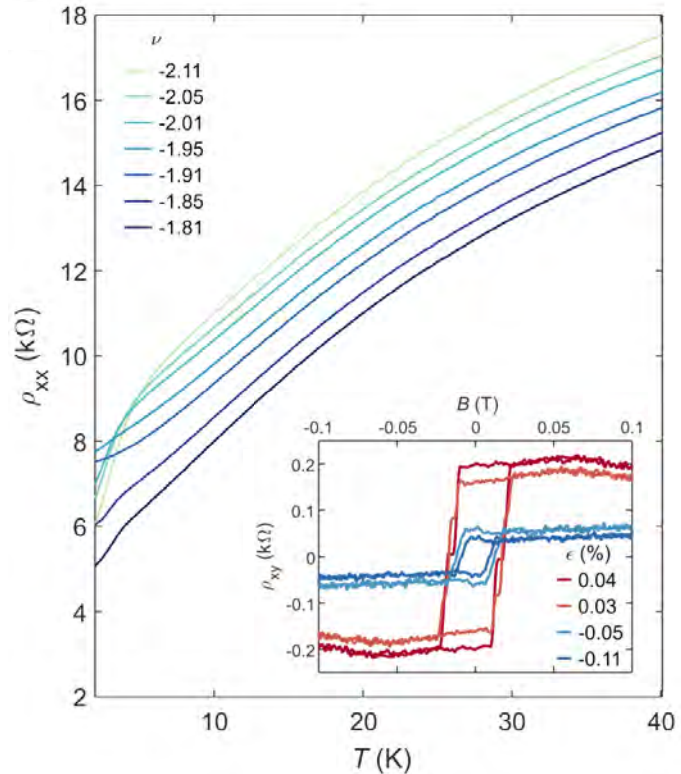
strain-induced band structure deformations may modify hybridization between light and heavy sectors

Possible signature of coherence-incoherence crossover

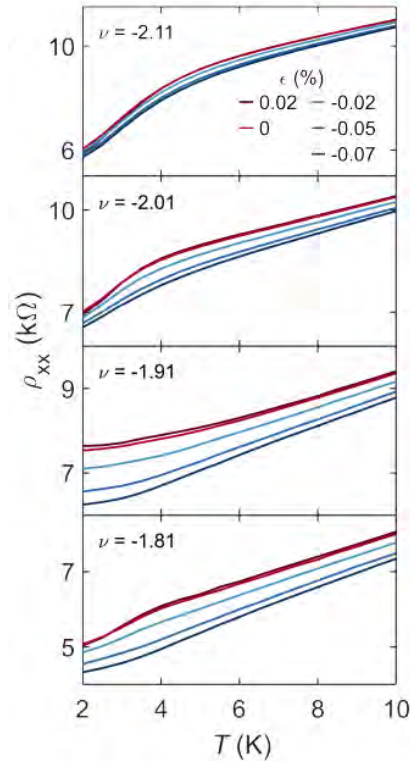
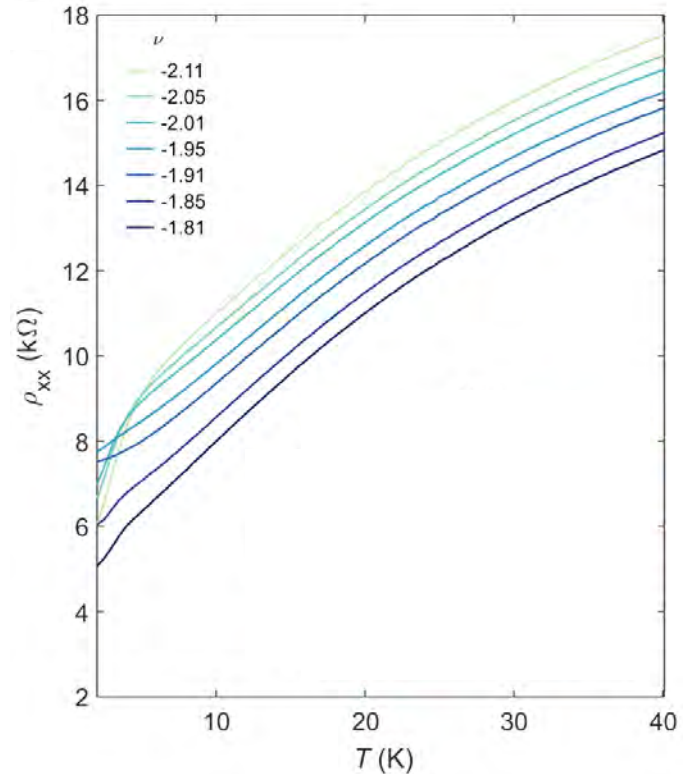


Diverging elastoresistance and ordered state formation

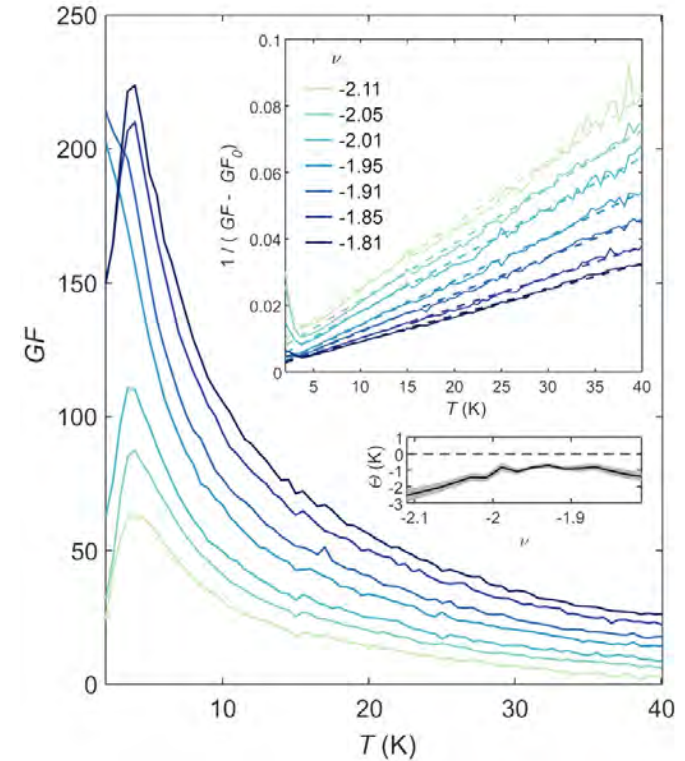
cutoff of the diverging elastoresistance associated with the formation of an isopin-ordered phase



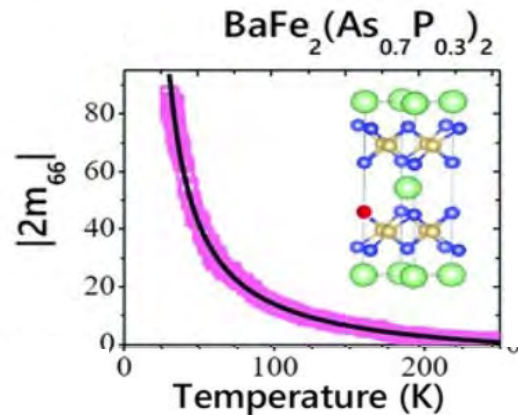
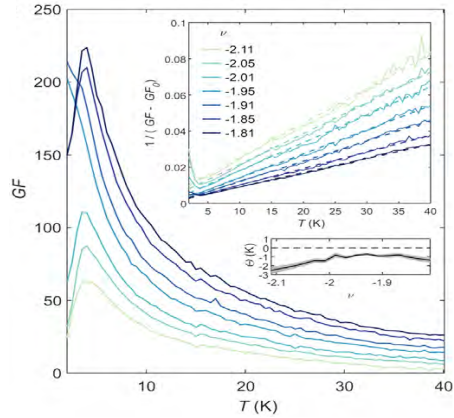
Curie-Weiss divergence of the elastoresistance



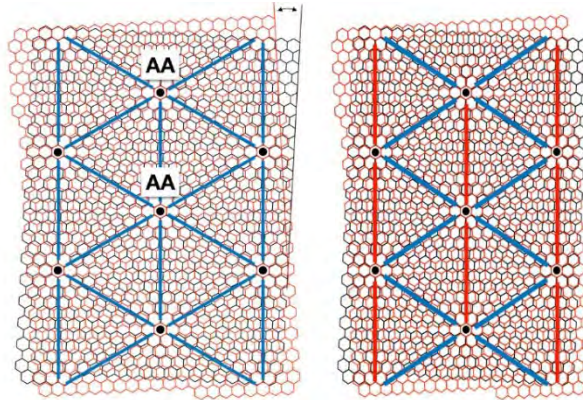
Curie-Weiss Law: $GF = \frac{C}{T - \Theta} + GF_0$



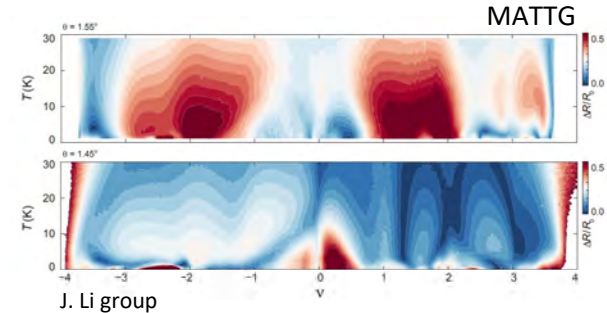
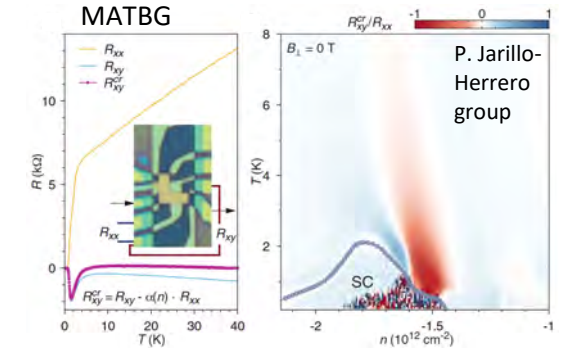
Possible role of nematic quantum critical fluctuations



nematic fluctuations



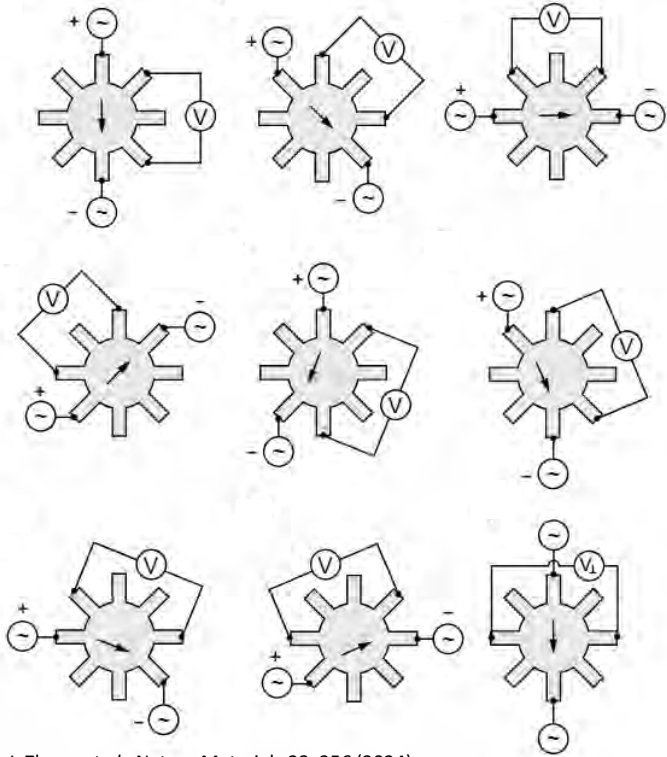
R. Fernandes & J. Venderbos, Science Advances 6, 8834 (2020)



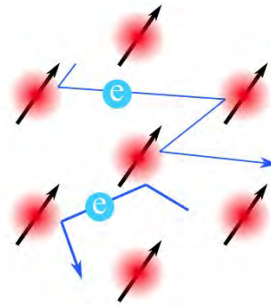
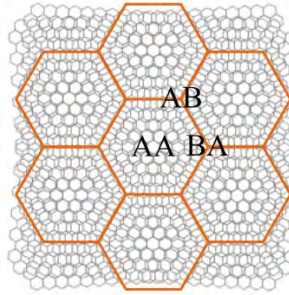
transport anisotropies hint at nematic ordering in MATBG (and MATTG)

Future work: symmetry-decomposed elastoresistance

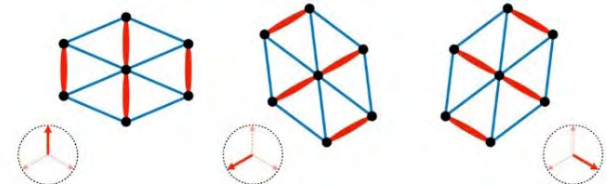
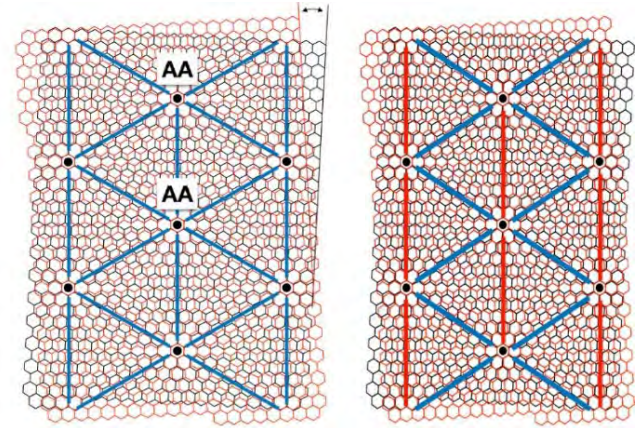
“sunflower” geometry for full resistivity tensor decomposition



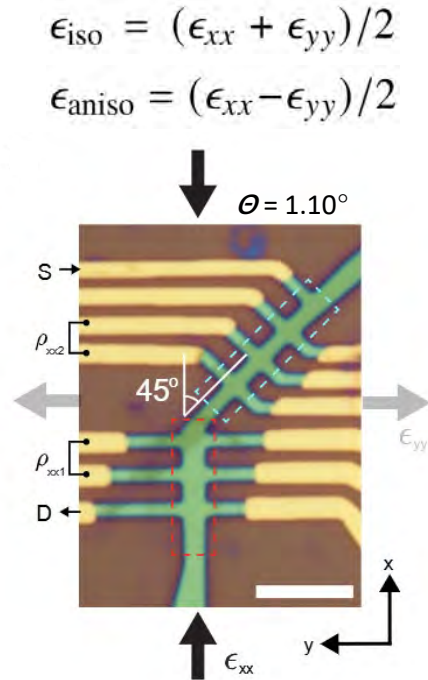
isotropic



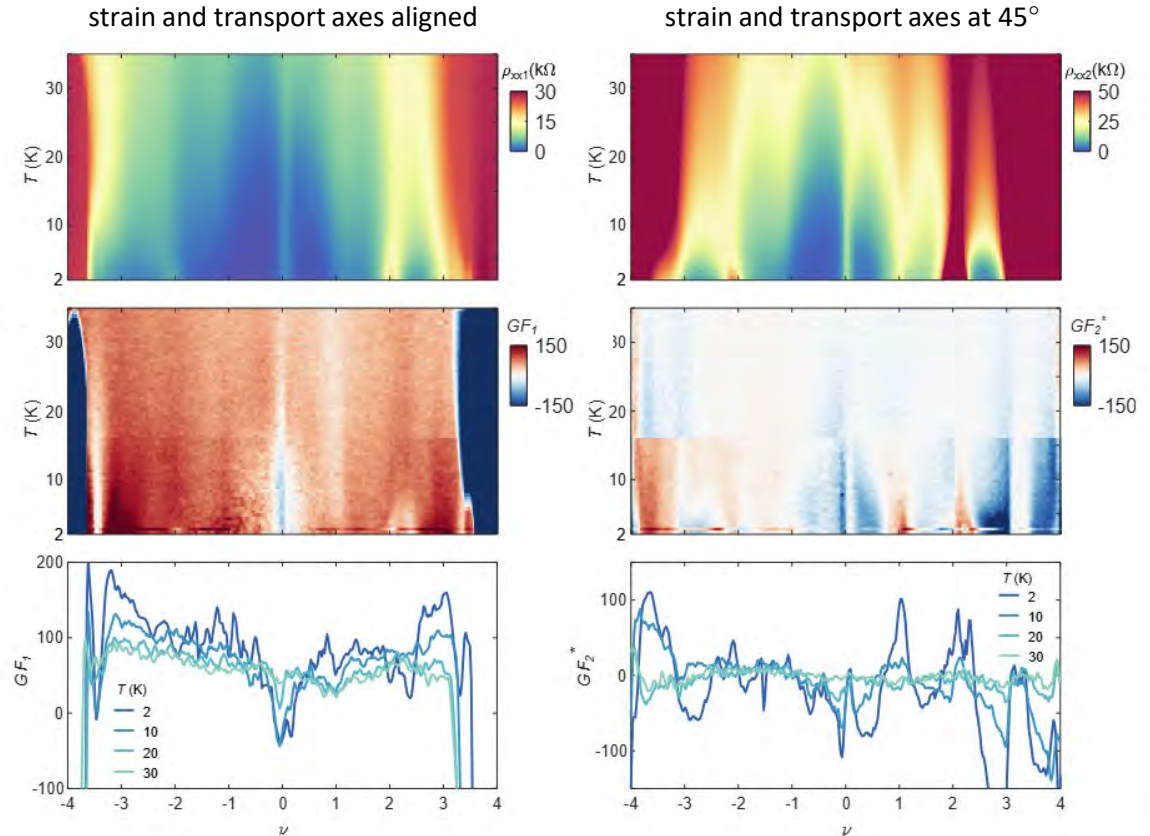
anisotropic



Preliminary attempt at symmetry resolution

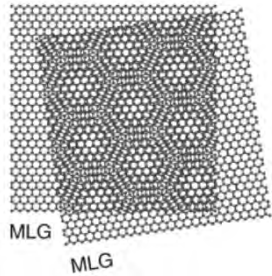


distinct strain response in two regions
implies contributions from both
symmetric and anisotropic channels

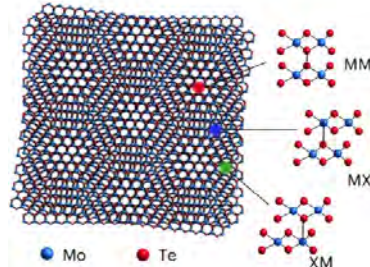


Strain for controlling van der Waals materials

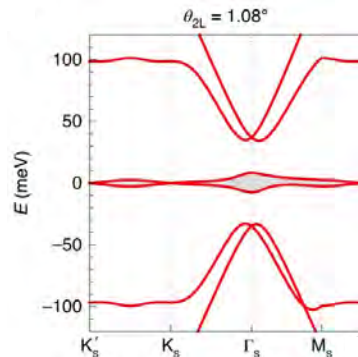
moiré multilayers



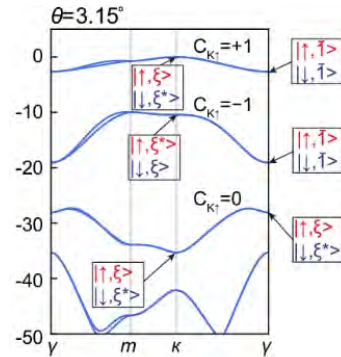
L. Balents et al., Nat. Phys. (2020)



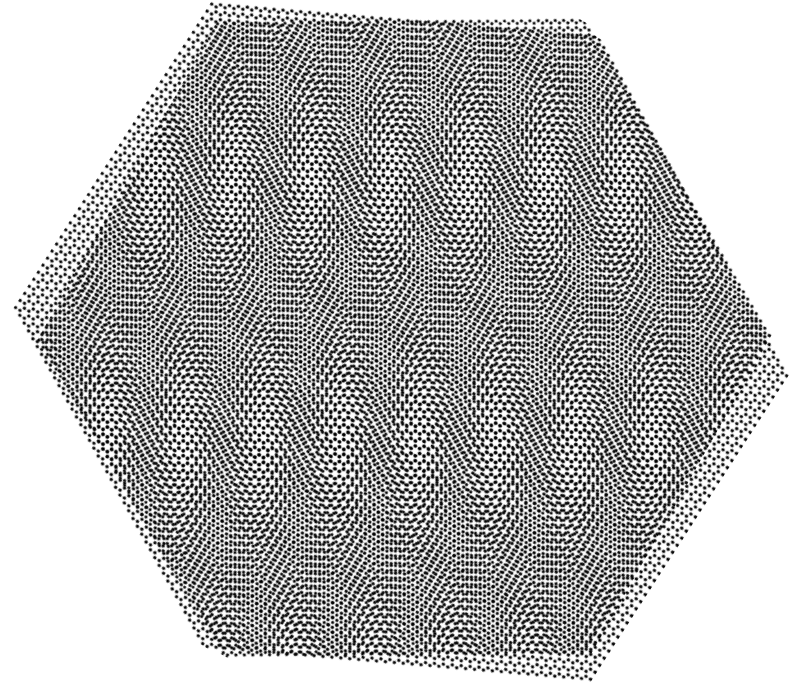
E. Thompson et al., Nat. Phys. (2025)



J. M. Park et al., Nat. Materials 21, 877 (2022)



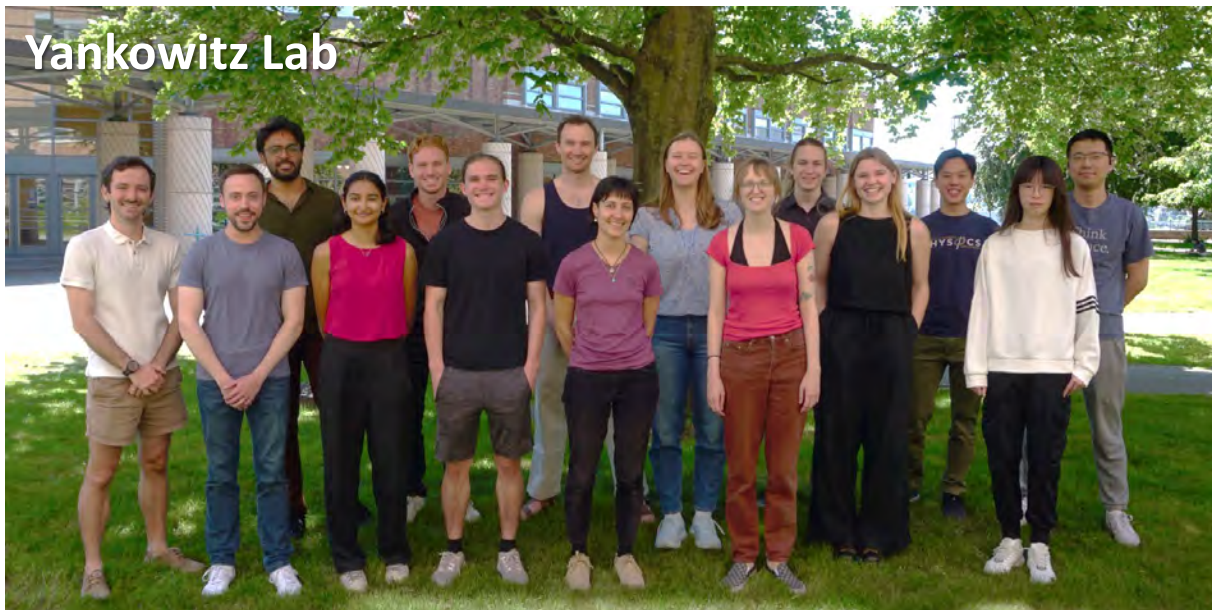
X.-W. Zhang et al., Nat. Comms. 15, 4223 (2024)



...toward new discoveries
with designer Hamiltonians

Acknowledgements

Yankowitz Lab



Back row (left to right):

Manish Kumar, Derek Waleffe, Eric Maginnis, Anna Okounkova, Isaac Mastel, Toby Chu, Xuetao Ma

Front row (left to right):

Jackson Chapman, Matthew Yankowitz, Aryana Bhattacharyya, Elijah Burns, Florie Mesple, Ellis Thompson, Abigail Soh, Violet Chen

Not pictured

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