

Friendly introduction to AdS/CMT

2. Non-equilibrium physics

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Retarded Green's functions

- ▶ consider a quantum system in thermal equilibrium:

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- ▶ then to first order in h ,

$$\langle \mathcal{A}(\mathbf{x}, t) \rangle = \int_{-\infty}^{\infty} ds G_{\mathcal{A}\mathcal{O}}^{\text{R}}(\mathbf{x}, \mathbf{y}, t - s) h(\mathbf{y}, s)$$

$$G_{\mathcal{A}\mathcal{O}}^{\text{R}}(\mathbf{x}, \mathbf{y}, t) = i\Theta(t) \langle [\mathcal{A}(\mathbf{x}, t), \mathcal{O}(\mathbf{y})] \rangle$$

Quasinormal modes dominate response

- ▶ singularity structure of correlators dominates response; e.g.

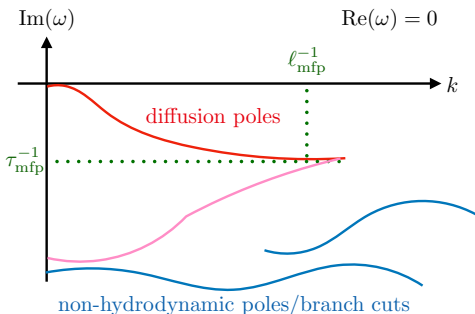
$$G^{\text{R}}(k, \omega) \sim \underbrace{\frac{1}{\omega + iDk^2}}_{\text{hydrodynamic diffusion}} + \underbrace{\frac{1}{\omega + i\tau^{-1}}}_{\text{finite lifetime}} + \dots$$

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- quasiparticles replaced by these quasnormal modes



Linear response in holography

- ▶ recall in holography, bulk field ϕ obeys

$$\phi = \underbrace{hr^{d+1-\Delta}}_{\text{source}} + \underbrace{\frac{\langle \mathcal{O} \rangle}{2\Delta - d - 1}}_{\text{response}} r^\Delta + \dots$$

with r the bulk radial coordinate

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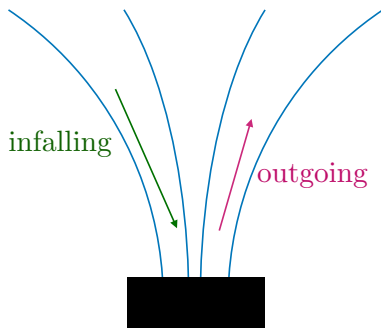
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- ▶ can generalize to evaluate higher-point (beyond linear response) and far from equilibrium

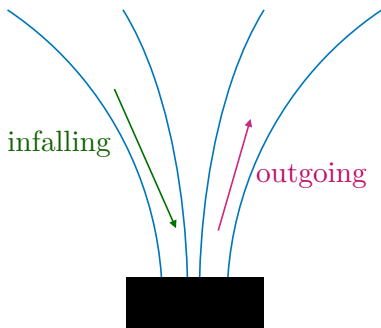
Infalling boundary conditions

- ▶ calculating $G^R \implies$ infalling boundary conditions in the presence of a black hole!



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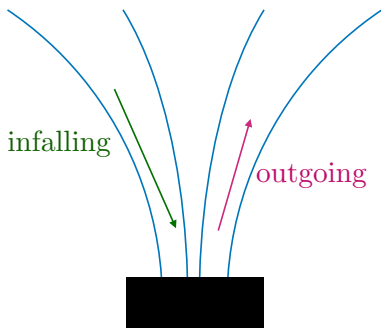
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- ▶ e.g. scalar field ϕ obeys

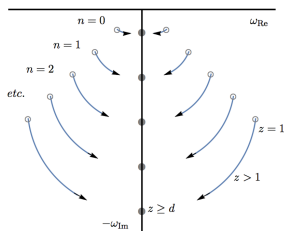
$$\phi(k, \omega, r \rightarrow r_h) \sim (r_h - r)^{-i\omega/4\pi T}$$

Holographic prescription for quasinormal modes

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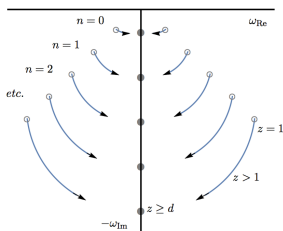
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- ▶ claim: every gapless holographic model has quasinormal mode at frequency ω_* with

$$\text{Im}(\omega_*) \gtrsim -T$$

(Planckian decay rate)

Diffusion

- ▶ a particular quasinormal mode: hydrodynamic diffusion

$$\partial_t \rho = D_{\text{charge}} \nabla^2 \rho = \frac{\sigma}{\chi_{\rho\rho}} \nabla^2 \rho,$$

$$\partial_t \epsilon = D_{\text{energy}} \nabla^2 \epsilon = \frac{T\kappa}{\chi_{\epsilon\epsilon}} \nabla^2 \epsilon,$$

$$\partial_t P_i = D_{\text{momentum}} \nabla^2 P_i = \frac{\eta}{\chi_{PP}} \nabla^2 P_i$$

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[Hartnoll; 1405.3651]

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[Hartnoll; 1405.3651]

- ▶ transport coefficients σ , κ , η fixed by physics on the horizon! (membrane paradigm)

[Iqbal, Liu; 0809.3808]

Viscosity bound

- ▶ all holographic models with Einstein gravity + matter have universal shear viscosity η : [Kovtun, Son, Starinets; hep-th/0405231]

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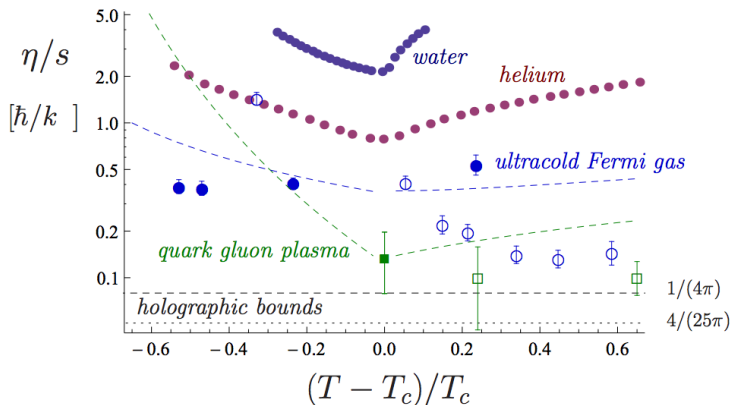
- ▶ in a charge neutral system, $\chi_{PP} = Ts$:

$$D_{\text{momentum}} \sim \frac{\eta}{Ts} \sim \frac{\hbar}{k_B T}$$

but at finite density, $\chi_{PP} \neq Ts...$

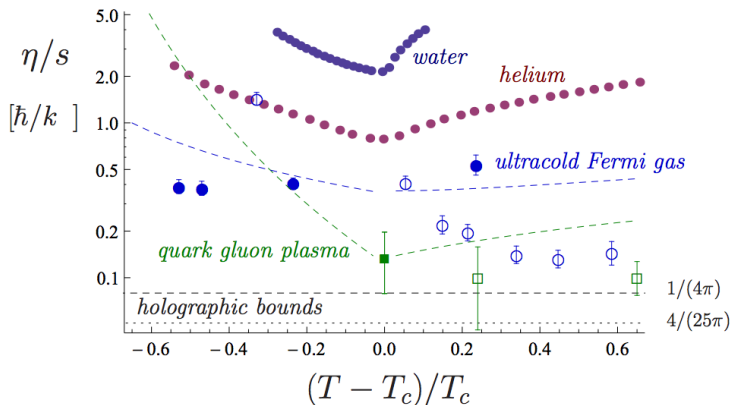
Viscosity bound in experiment

- bound consistent with experiment: [Adams *et al*; 1205.5180]



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- theoretically, bound has been violated [Brigante *et al*; 0712.0805]

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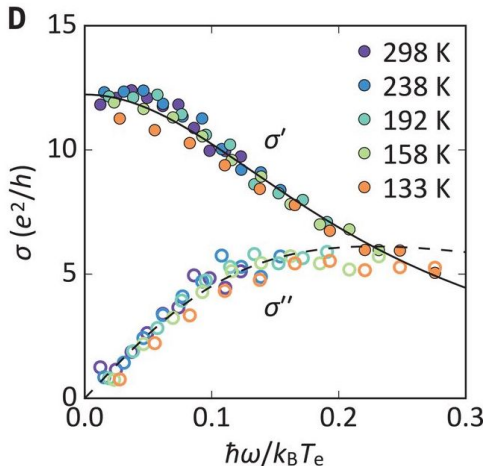
- ▶ calculating F is *very hard*:
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- ▶ holographic calculation of F : ODE in Mathematica

$$\partial_r (Y_1(r) \partial_r A_x) = -\omega^2 Y_2(r) A_x$$

with $Y_{1,2}$ known from bulk geometry

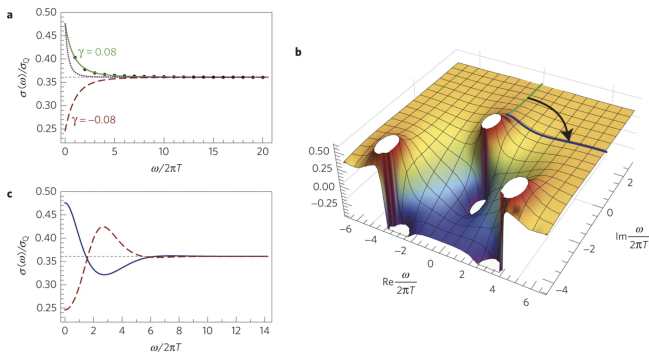
Conductivity at zero density: graphene

- recent experiment on graphene measured $\sigma = F(\omega/T)$:



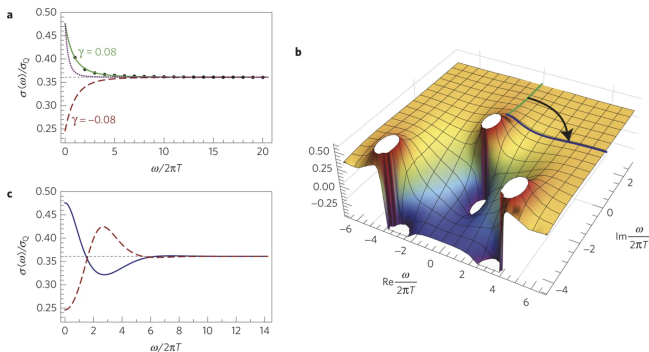
Analytic continuation

- ▶ using holography to analytically continue F in a 2d lattice model: [Witczak-Krempa *et al*; 1309.2941]



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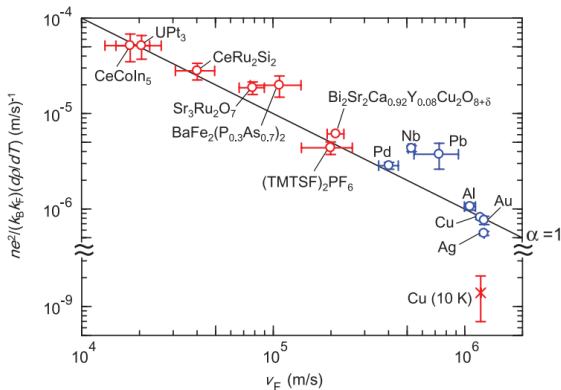
- ▶ note: $\omega \gg T$ region also understood with conformal perturbation theory [Lucas *et al*; 1608.02586]

Planckian time in experiment?

- Planckian resistivity observed in many strange metals:

$$\rho \sim \frac{m}{ne^2} \frac{k_B T}{\hbar}$$

[Bruin *et al.*; (2013)]



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- ▶ σ sensitive to how translational symmetry broken

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does this really work? what is τ ?

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- ▶ holography reproduces this result!

[Lucas; 1501.05656]

$$\lim_{\omega \rightarrow 0} \frac{\text{Im}(G_{\mathcal{O}\mathcal{O}}^R(\mathbf{k}, \omega))}{\omega} \sim \sqrt{g_{\text{horizon}}} \phi(r_h)^2$$

because horizon physics determines spectral weight!

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- ▶ near spin density wave criticality? [Patel, Sachdev; 1408.6549]

$$\rho \sim T^0 + T + T^4$$

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with $\Phi = mx$ (an exact solution) [Andrade, Withers; 1311.5157]

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- ▶ field theory analogues in U(1)-symmetric theories?

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- ▶ e.g. AdS-Einstein-Maxwell (minimal) model, $d = 2$:

$\sigma \geq 1$ (conductivity of clean neutral plasma)

$$\kappa \geq \frac{4\pi^2 T}{3}$$

[Grozdánov, Lucas, Sachdev, Schalm; 1507.00003]

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- ▶ bounds saturated by mean field model with $m = \infty$

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- ▶ the usual “metal-insulator” transition is

$$C_1(T) \sim T^\alpha, \quad \alpha \text{ tunable through } 0$$

and is neither Anderson or Mott

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with $C_{1,2}(T) \sim T^{\alpha_{1,2}}$ having tunable scaling

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- ▶ many mean field models are unstable at large m – is there a sensible endpoint?

Diffusion and chaos

- ▶ some universality: the *thermal diffusion* constant obeys

$$D_{\text{thermal}} = c \frac{v_{\text{B}}^2}{T}$$

for $O(1)$ constant c , in homogeneous systems

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- ▶ this relation fails in inhomogeneous systems:

$$D_{\text{thermal}} \leq c \frac{v_{\text{B}}^2}{T}$$

(left hand side can be arbitrarily smaller) [Gu *et al*; 1702.08462]

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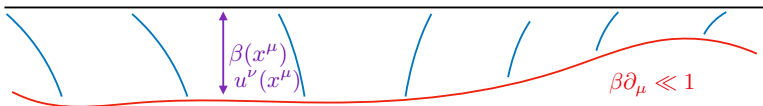
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- ▶ but ... numerical general relativity = hard!

Fluid-gravity correspondence

- ▶ long wavelength dynamics of black holes reproduces nonlinear hydrodynamics!

$$ds^2 = \frac{L^2}{r^2} [-2dr u^\mu(x) dx_\mu + (1 - f(r, \beta(x)))(u^\mu(x) dx_\mu)^2 + dx^\mu dx_\mu]$$

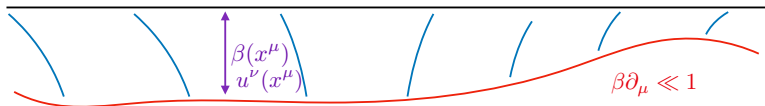


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- ▶ this method can be used to generate high order corrections to fluid dynamics: e.g.

$$\partial_t \beta = D \nabla^2 \beta + D_4 \nabla^4 \beta + D_6 \nabla^6 \beta$$

Borel resummability of hydrodynamics can be investigated

[Grozdanov, Kovtun, Starinets, Tadic; 1904.01018]

Quantum quenches

- ▶ consider a rapid change in the Hamiltonian:

$$H = H_{\text{CFT}} + \lambda \text{sech} \frac{t}{\tau} \mathcal{O}$$

with \mathcal{O} an operator of dimension Δ

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- ▶ (numerical) correlators tractable in these time-dependent backgrounds

Open questions

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- ▶ questions raised by AdS/CMT:
 - ▶ thermal chaos/diffusion in spin chains?
 - ▶ “incoherent” metals? resistivity saturation at strong coupling?
 - ▶ CFTs at finite T – defects? OPE?